

UC Davis High Energy Seminar

Ultrasensitive Searches for the Axion

Karl van Bibber, LLNL

January 23, 2007

AXION

Outline

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Some basics about the axion

Three experimental fronts:

Searches for halo dark matter

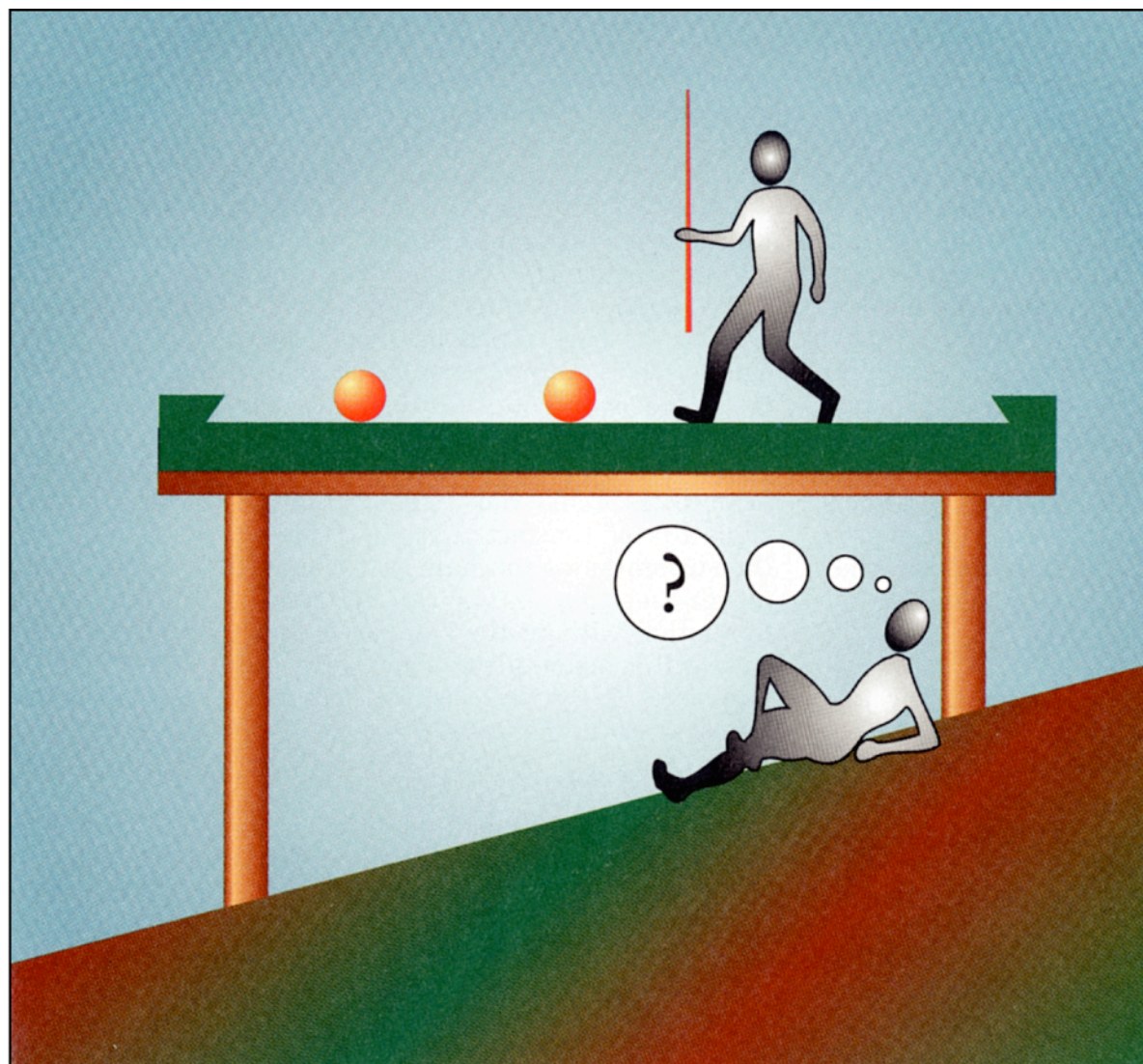
Searches for solar axions

Purely laboratory experiments

Final remarks

TSP's* fine-tuning problem

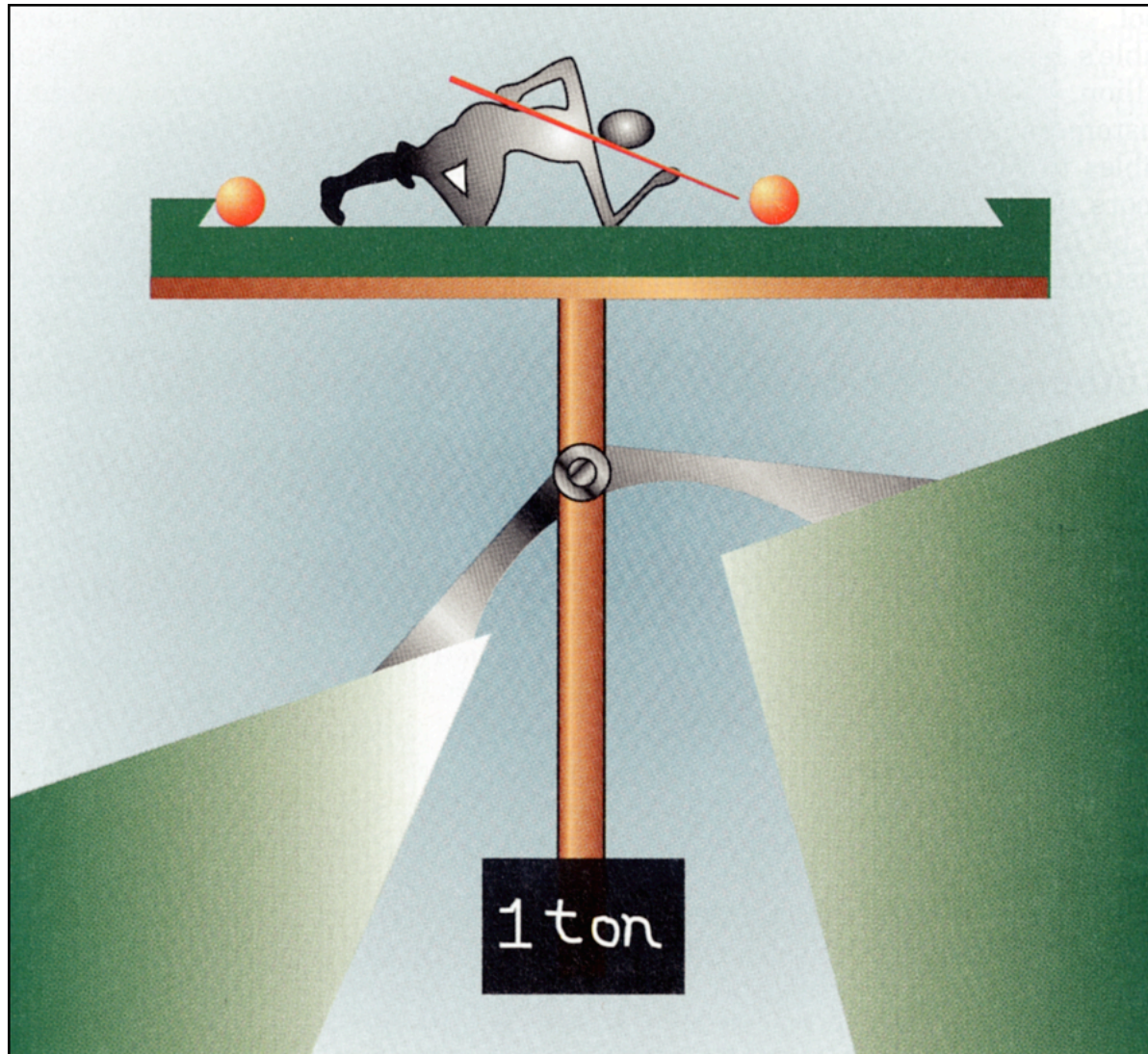
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*Thinking Snookers Player (Pierre Sikivie, Physics Today 49 (1996)22)

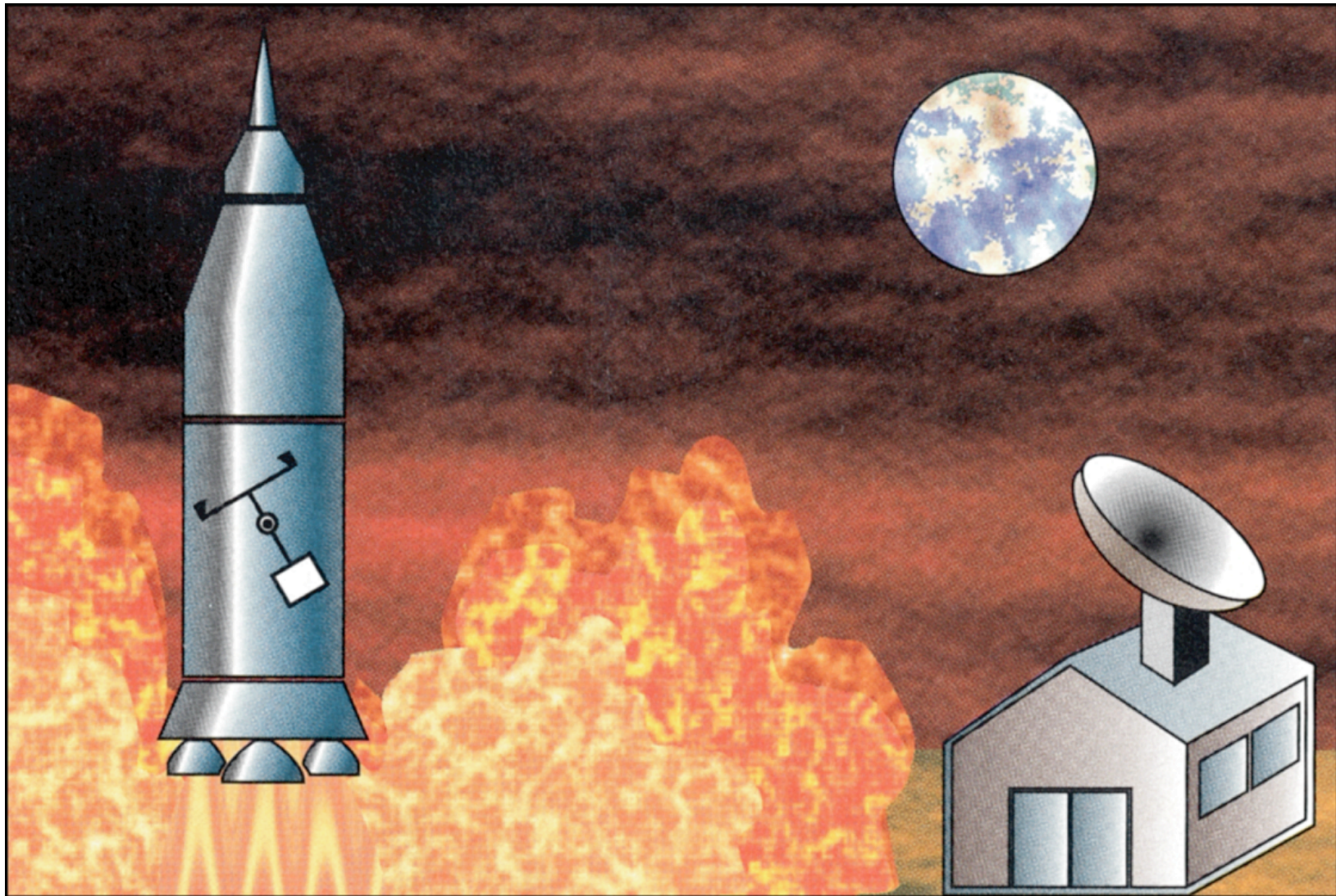
TSP's hypothesis, and first unsuccessful experiment

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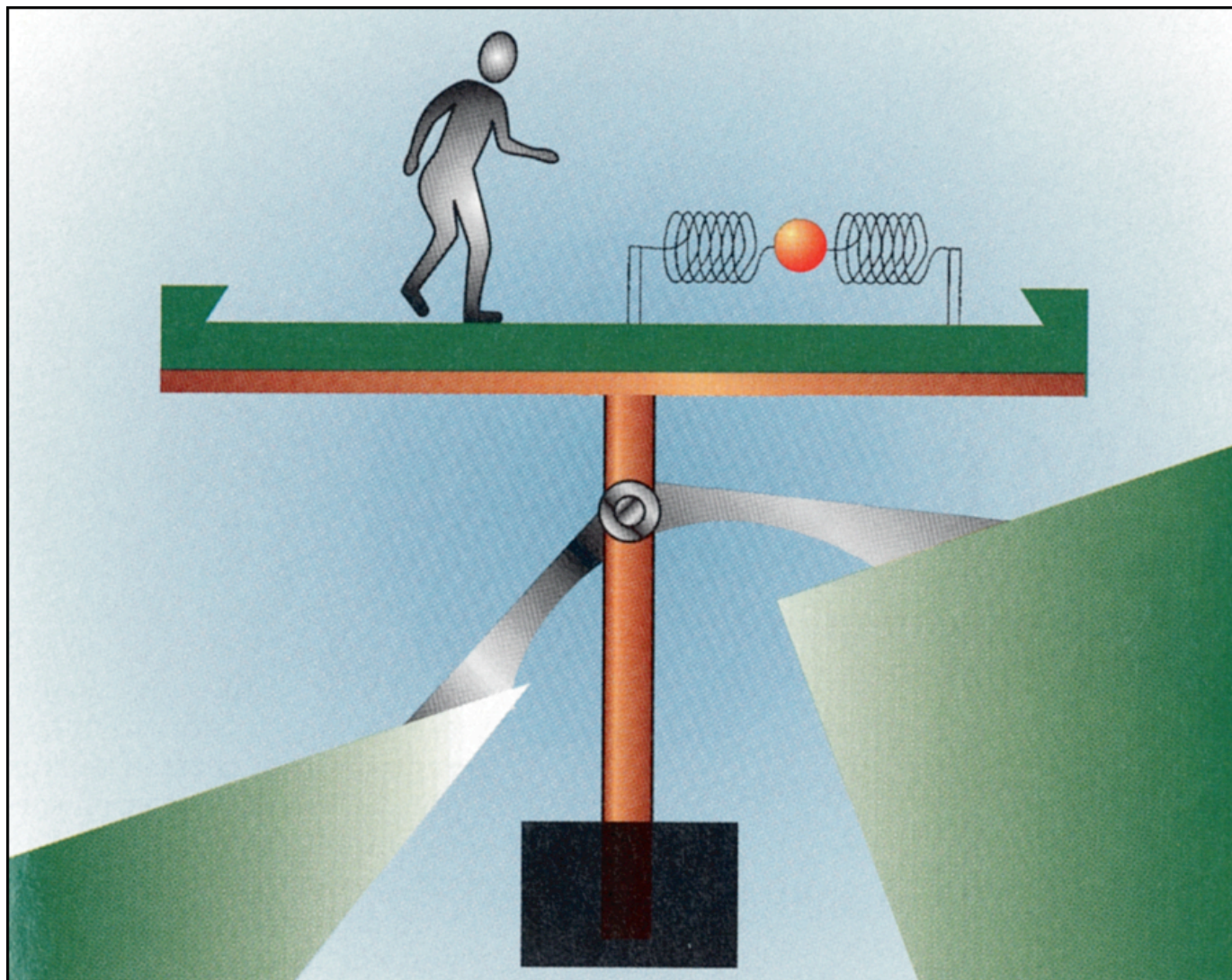
The key insight

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A high-Q search for relic oscillations

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The Axion

The Strong-CP Problem

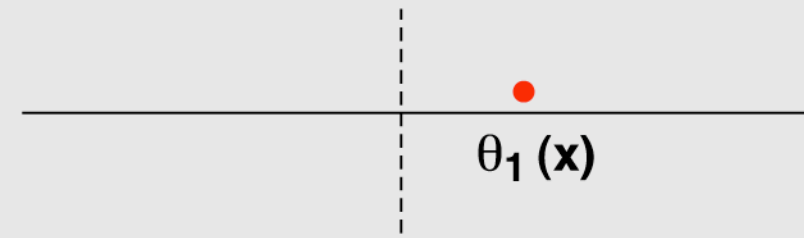
- $\mathcal{L}_{\text{QCD}} = \dots + \frac{\theta}{32\pi^2} \mathbf{G}\tilde{\mathbf{G}}$
 - Explicitly CP-violating
- But neutron e.d.m. $|d_n| < 10^{-25} \text{ e} \cdot \text{cm}$
 - $\bar{\theta} < 10^{-10}$
 - Strong-CP preserving

$$\text{CP} \left(\begin{array}{c} \uparrow \mu_n \\ \uparrow d_n \\ \text{In} \rangle \\ \downarrow \\ \downarrow \mu_n \end{array} \right) = \begin{array}{c} \uparrow d_n \\ \downarrow \\ \downarrow \mu_n \end{array} \neq \text{In} \rangle$$

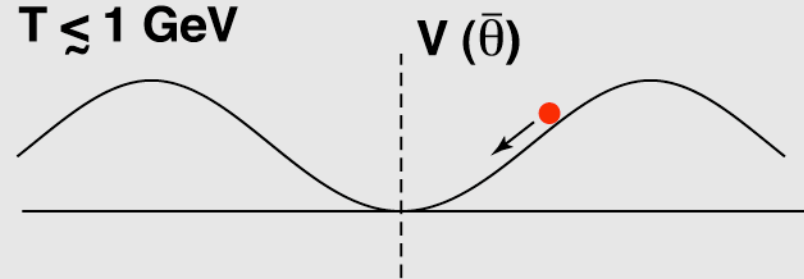
- Why?

Peccei-Quinn / Weinberg-Wilczek

- θ a dynamical variable
- $T = f_a$ spontaneous symmetry breaking



- $T \lesssim 1 \text{ GeV}$



- $\bar{\theta}$ dynamically $\rightarrow 0$
- Remnant oscillation = Axion

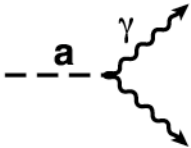
Completing the analogy $f \leftrightarrow l$

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	PQ-symmetry breaking scale	Pendulum length
Quanta m_a (w)	$\sim f^{-1}$	$\sim l^{-1/2}$
Couplings $g_{a ii}$	$\sim f^{-1}$	$\sim l^{-1}$
Total energy W_a (E)	$\sim f^{7/6}$	$\sim l$

Properties of the Axion

- The Axion is a light pseudoscalar resulting from the Peccei-Quinn mechanism to enforce strong-CP conservation
- f_a , the SSB scale of PQ-symmetry, is the one important parameter in the theory

<p>Mass and Couplings</p> $m_a \sim 6 \mu\text{eV} \cdot \left(\frac{10^{12} \text{ GeV}}{f_a} \right)$ <p>Generically, all couplings</p> $g_{a ii} \propto \frac{1}{f_a}$	<p>Cosmological Abundance</p> $\Omega_a \sim \left(\frac{5 \mu\text{eV}}{m_a} \right)^{7/6}$ <p>(Vacuum misalignment mechanism)</p>
<p>Coupling to Photons</p>  $g_{a\gamma\gamma} = \frac{\alpha g_\gamma}{\pi f_a}; g_\gamma = \begin{cases} 0.97 \text{ KSVZ} \\ -0.36 \text{ DFSZ} \end{cases}$	<p>Axion Mass 'Window'</p> $10^{-(5 \text{ to } 6)} \text{ eV} < m_a < 10^{-(2 \text{ to } 3)} \text{ eV}$ <p>(Overclosure) (SN1987a)</p> <p>With lower end of window preferred if $\Omega_{\text{CDM}} \sim 1$</p>

Microwave cavity searches for axionic dark matter *AXION*

Some basics about dark matter

Principle of the Sikivie experiment

The first generation experiments (RBF, UF) c. 1990

Axion Dark Matter eXperiment (ADMX) @ LLNL

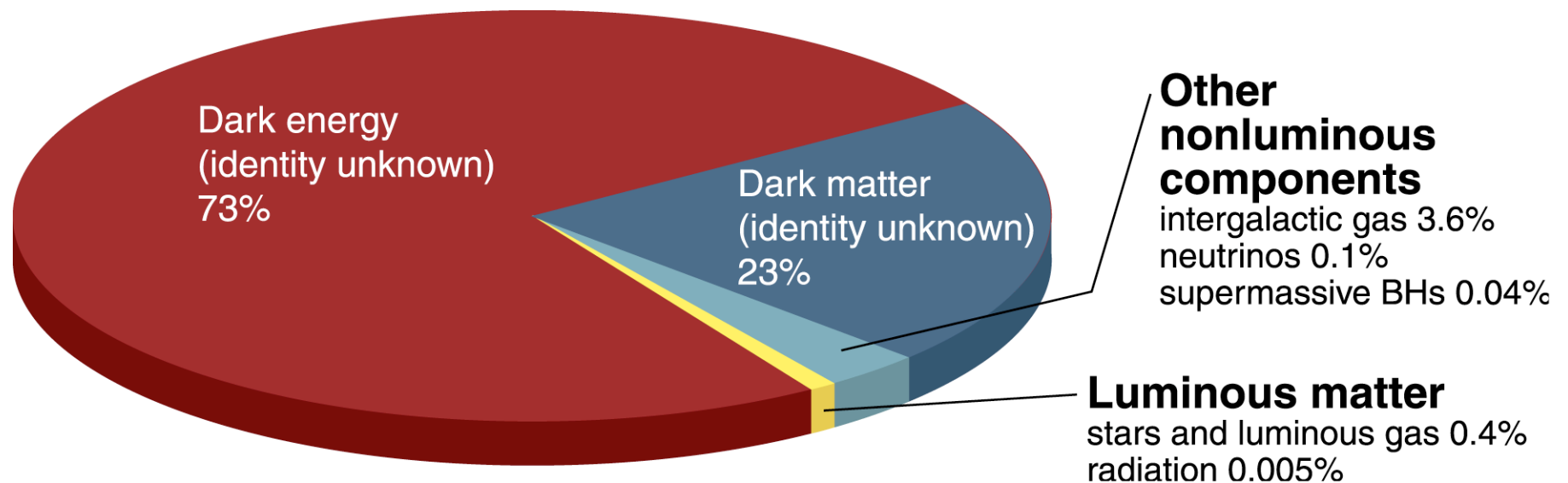
Upgrade based on quantum-limited SQUID amplifiers

The Rydberg-atom single-quantum detector @ Kyoto

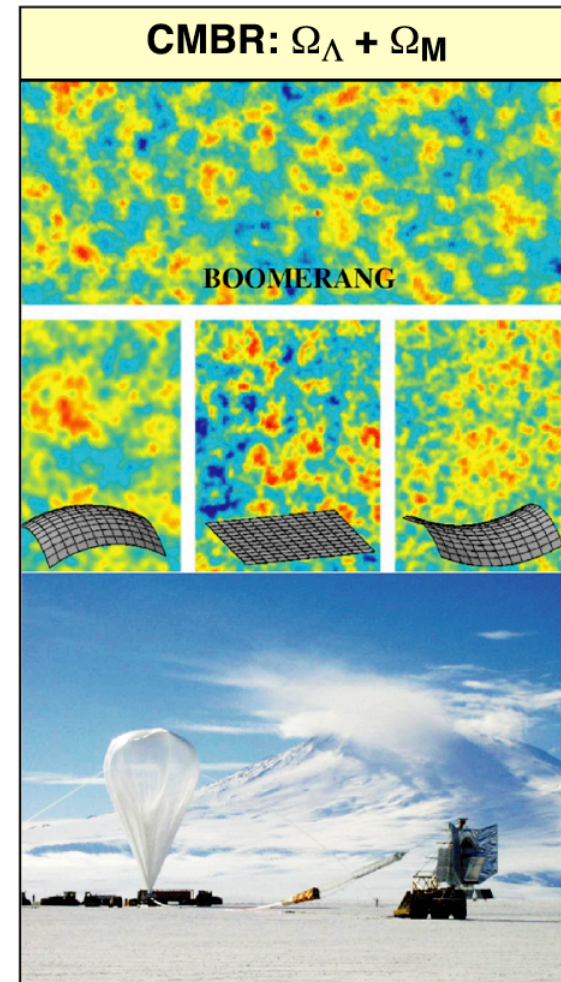
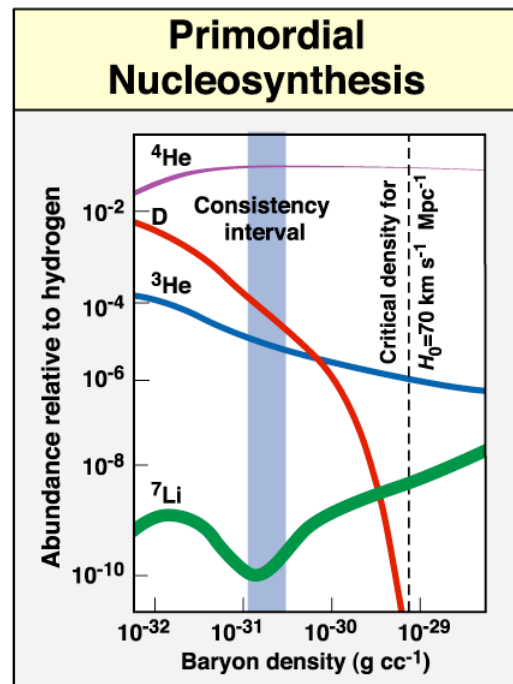
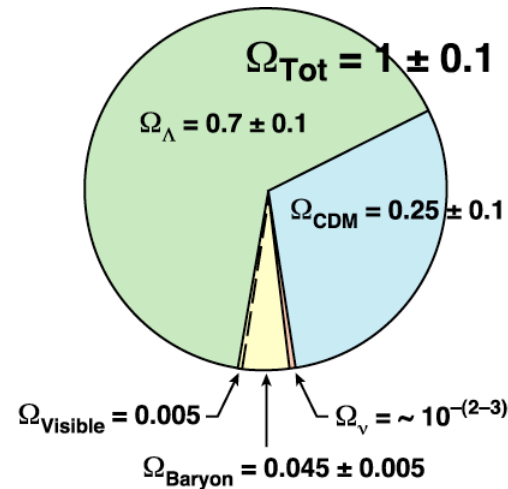
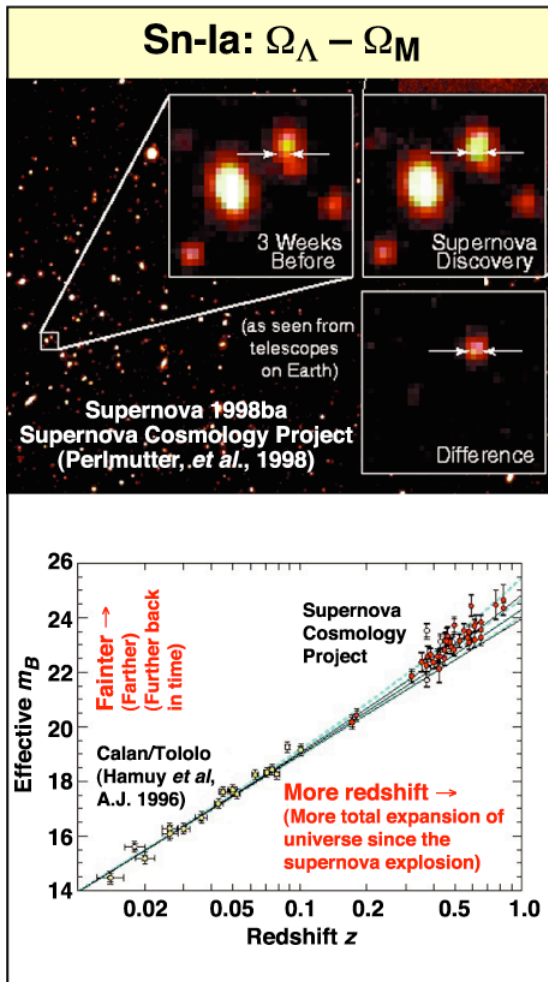
The cosmological inventory is now well-delineated

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- But we know neither what the “dark energy” or the “dark matter” is
- A particle relic from the Big Bang is strongly implied for DM
 - WIMPs ?
 - Axions ?

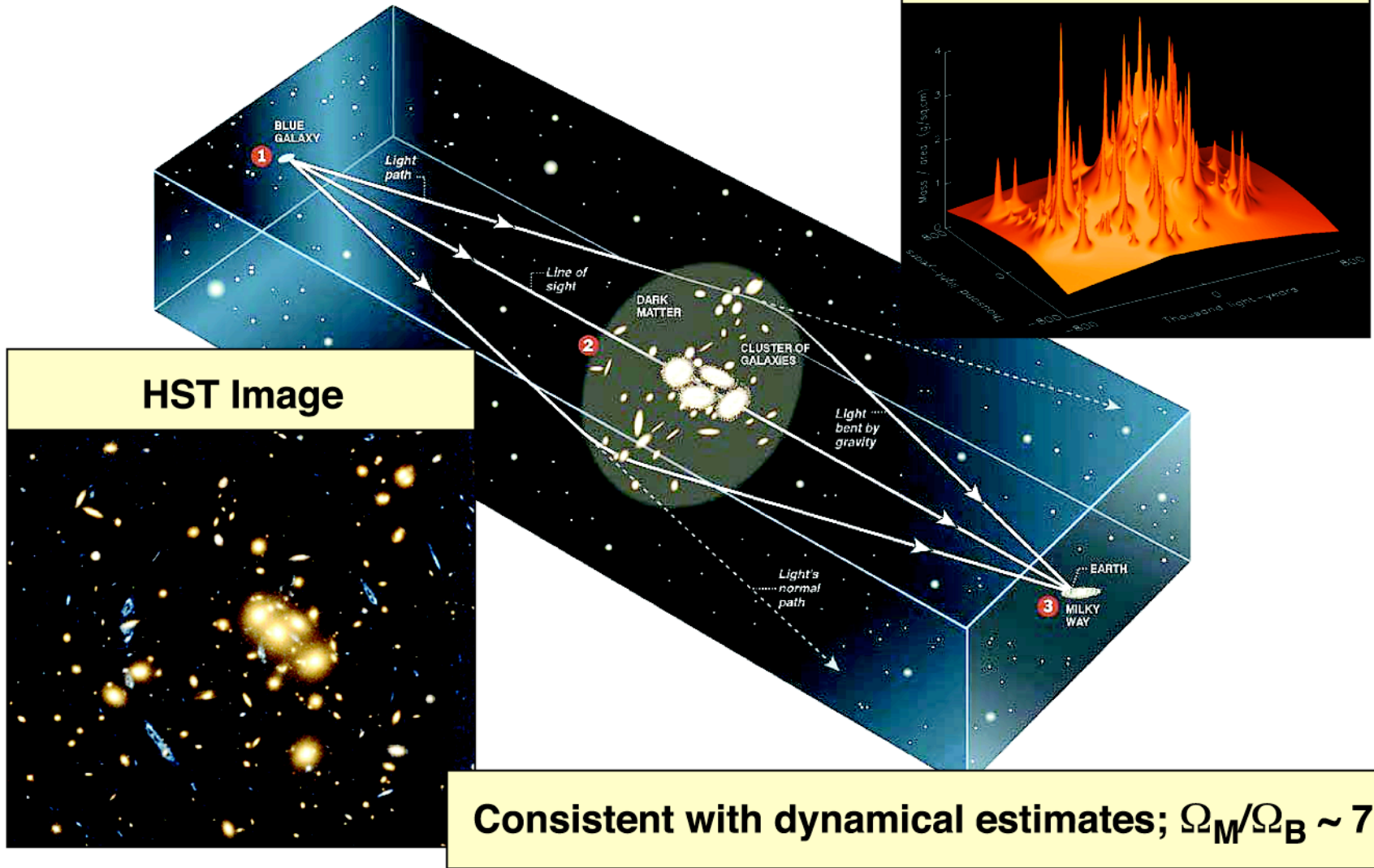


The advent of “precision cosmology”



Cluster lensing of background galaxy

- Cluster mass reconstruction from multiple gravitational lensing of background galaxy



Consistent with dynamical estimates; $\Omega_M/\Omega_B \sim 7$

Rotation Curves — Galactic Dark Matter

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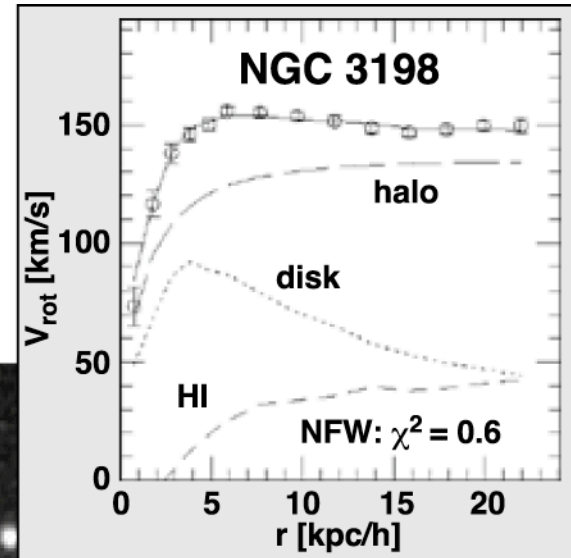
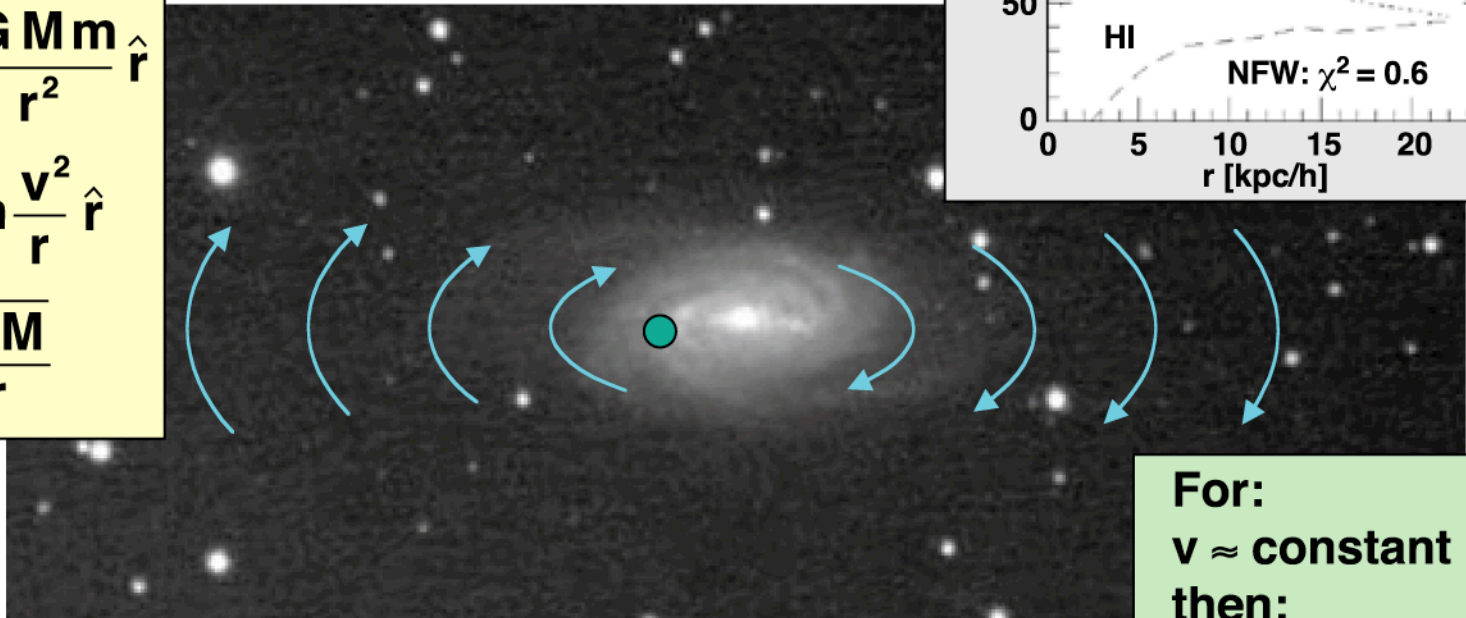
- Galaxies have constant rotation curves

If mass were localised

$$\mathbf{F} = -\frac{GMm}{r^2} \hat{\mathbf{r}}$$

$$= -m \frac{v^2}{r} \hat{\mathbf{r}}$$

$$v = \sqrt{\frac{GM}{r}}$$



- Dark Matter is right in our own Milky Way!
- Maximum likelihood local density $\rho \sim 450 \text{ MeV/cm}^3$

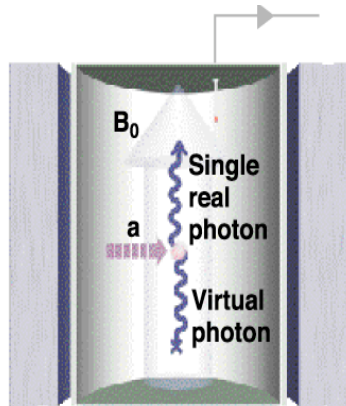
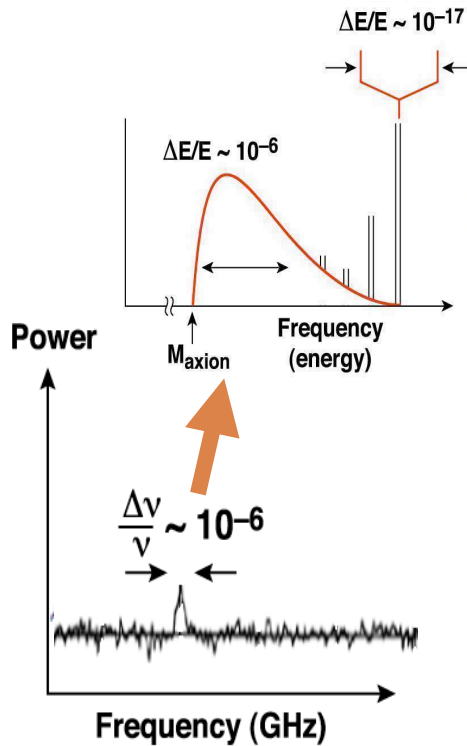
For:
 $v \approx \text{constant}$
then:

$$M(r) \propto r$$

$$M_{\text{dark}} \geq 10 M_{\text{lum}}$$

Nature of axionic dark matter, and principle of the microwave cavity experiment [Pierre Sikivie, PRL 51, 1415 (1983)]

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Axionic dark matter is very dense

Milky Way density: $\rho_{halo} \approx 450 \text{ MeV} \cdot \text{cm}^{-3}$
 Thus if $m_a \sim 10 \mu\text{eV}$: $\rho_{\#} \approx 10^{14} \text{ cm}^{-3}$

Axionic dark matter is highly coherent

$\beta_{virial} \approx 10^{-3} \rightarrow \lambda_{De\ Broglie} \approx 100 \text{ m}$
 $\Delta\beta_{flow} \approx 10^{-7} \rightarrow \lambda_{Coherence} \approx 1000 \text{ km}$

Resonance condition: $hn = m_a c^2 [1 + O(b^2 \sim 10^{-6})]$

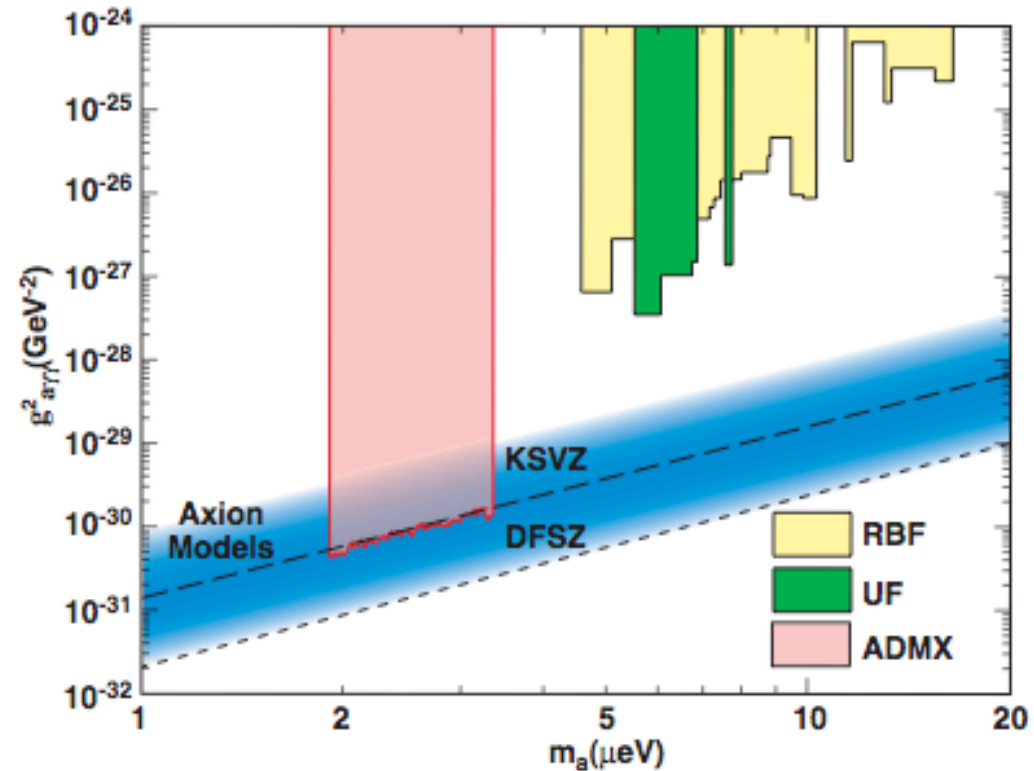
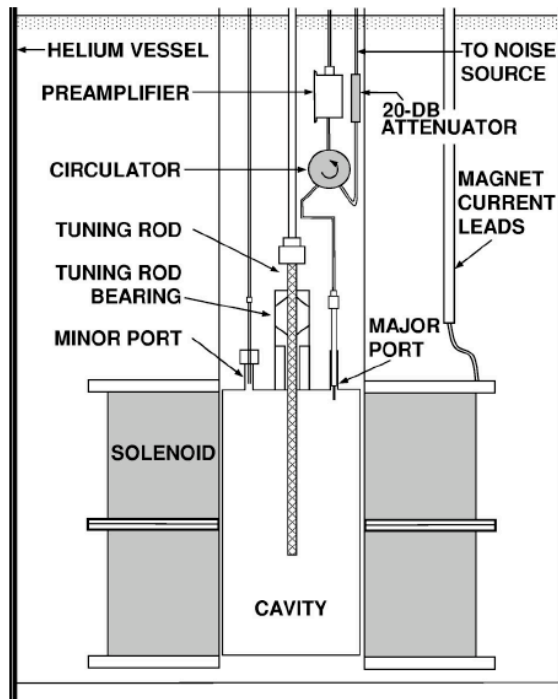
Signal power: $P \propto (B^2 V Q_{cav}) (g^2 m_a r_a) \sim 10^{-23} \text{ W}$

The microwave cavity experiment measures the *total energy* of the axion, thus revealing both Doppler motion and coherence of the axion fluid

The first-generation experiments RBF, UF – 1980's

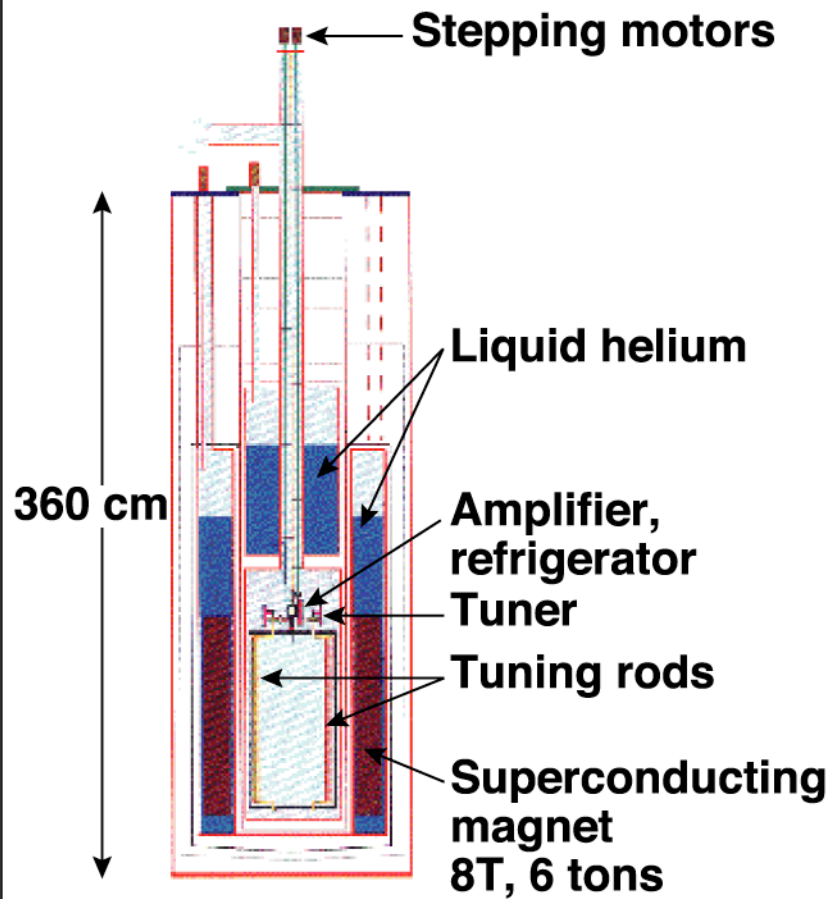
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From W. Wuensch *et al.*,
Phys. Rev. D40 (1989) 3153



The first-generation experiments already came within a factor of 100-1000 of the desired sensitivity – a stunning achievement

Magnet with Insert (side view)



Pumped LHe \rightarrow $T \sim 1.5$ k

Magnet (Wang NMR Inc.)

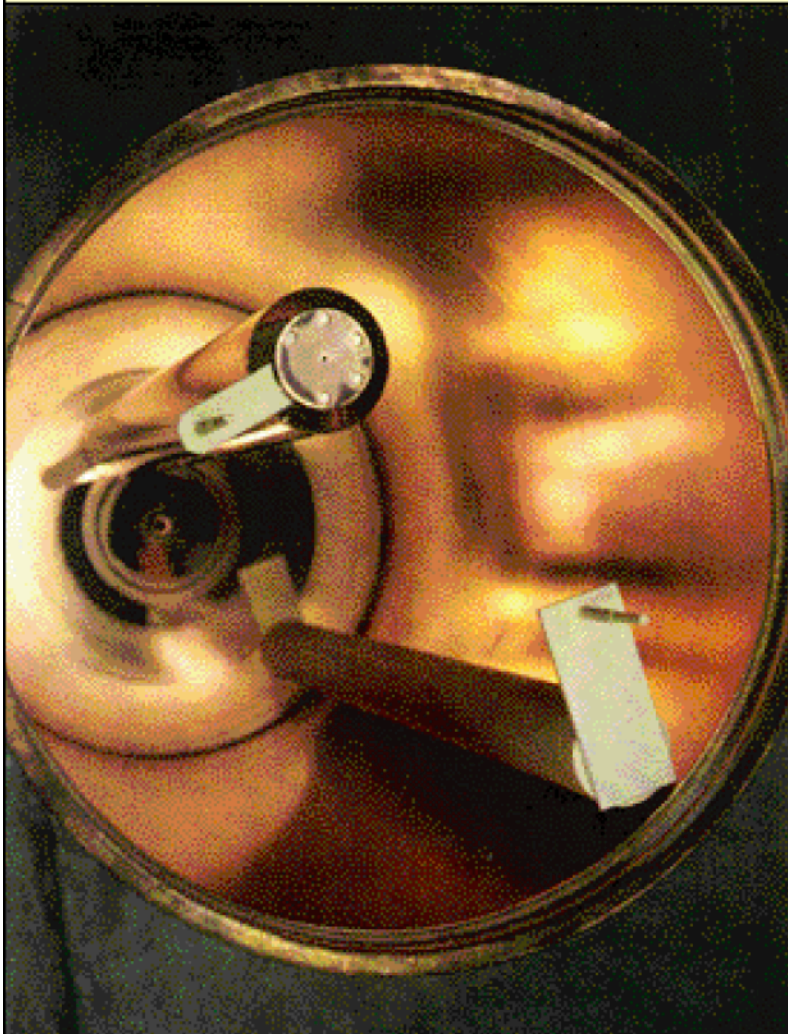


8 T, 1 m \times 60 cm \varnothing

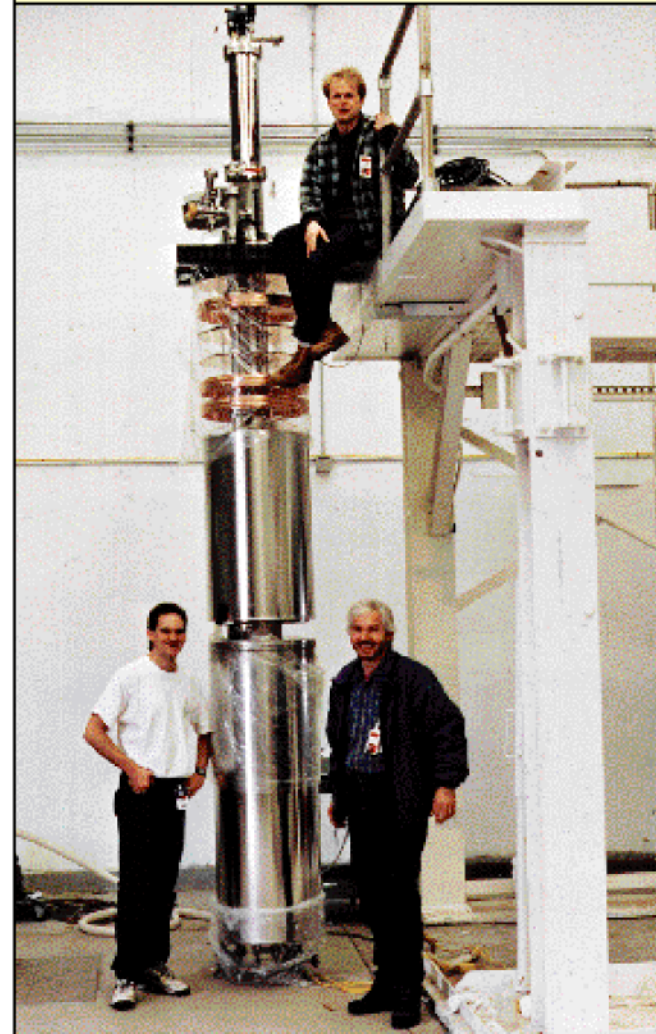
Axion hardware (cont'd)

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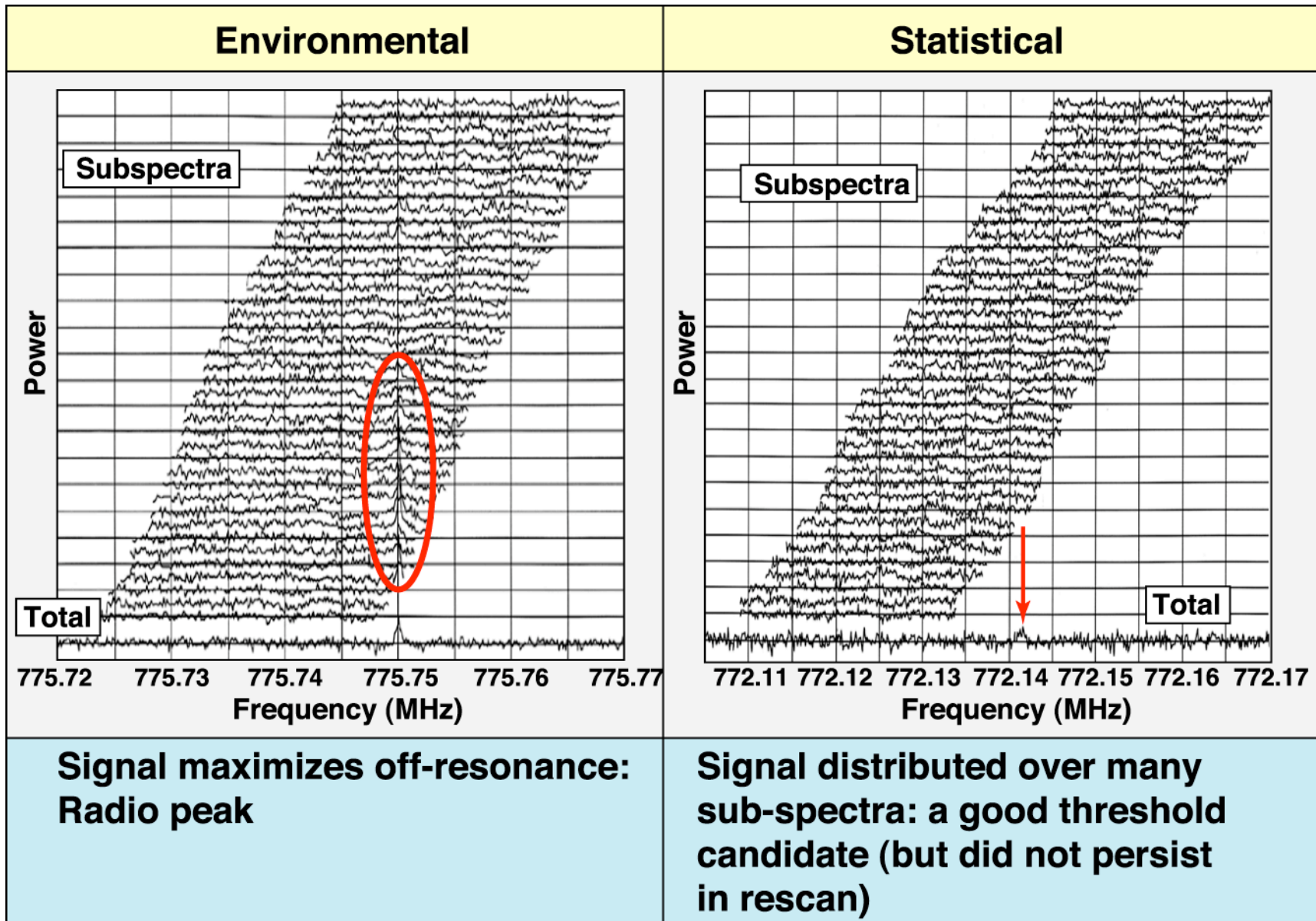
High-Q Cavity (~200,000)



Experimental Insert



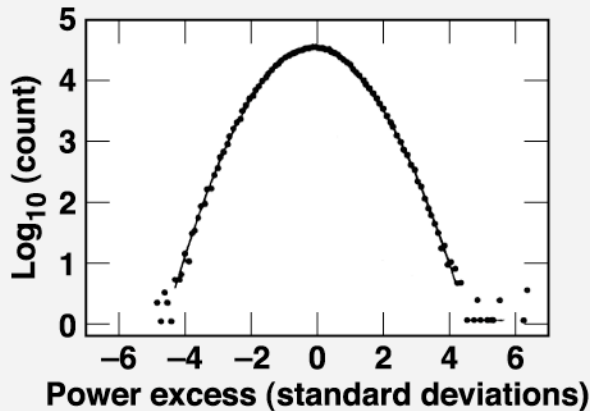
Sample data and candidates



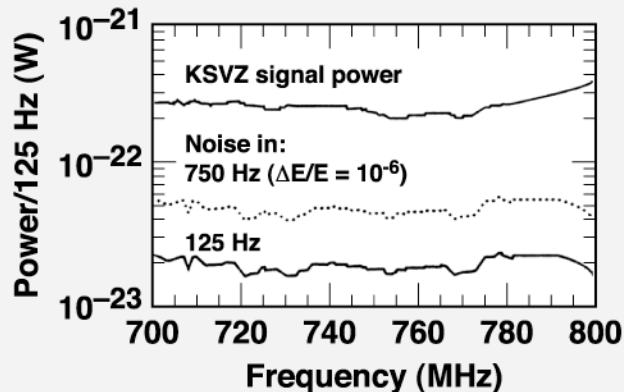
Brief outline of analysis – 100 MHz of data

Data, with Theoretical Curve

(Gaussian noise through receiver and analysis)



S/N > 4 for Thermalized KSVZ Axion

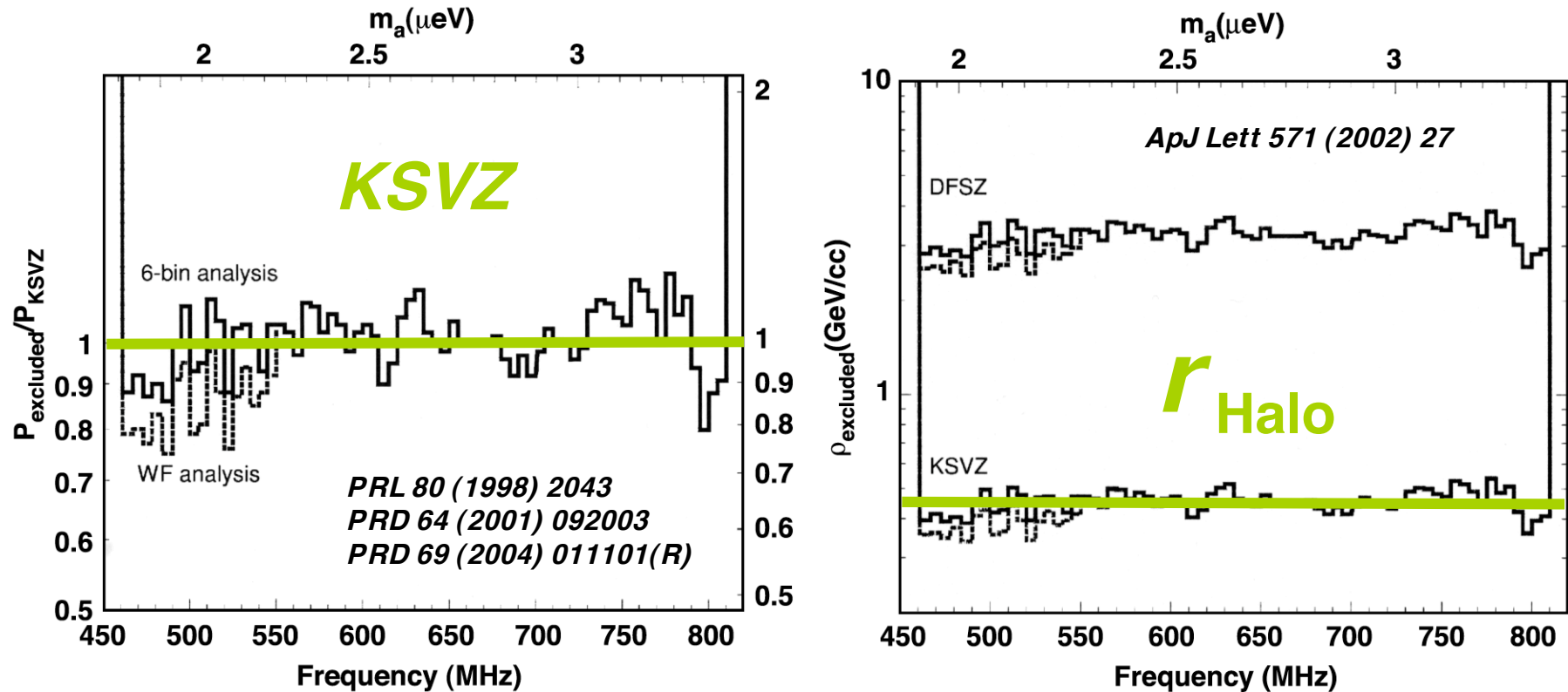


- Each frequency appears in >45 subspectra
- Weighted and co-added to produce spectrum
- 800,000 bins (125 Hz)/100 MHz

- 6535 candidates $> 2.25 \sqrt{6 \bar{\sigma}}$ (95% C.L.)
- Rescan all to same sensitivity
- 23 candidates (Net 90% C.L.)
- Each examined: radio peaks

For a persistent peak, the ultimate test is to turn off the magnet!

Limits on axion models and local axion halo density AXION

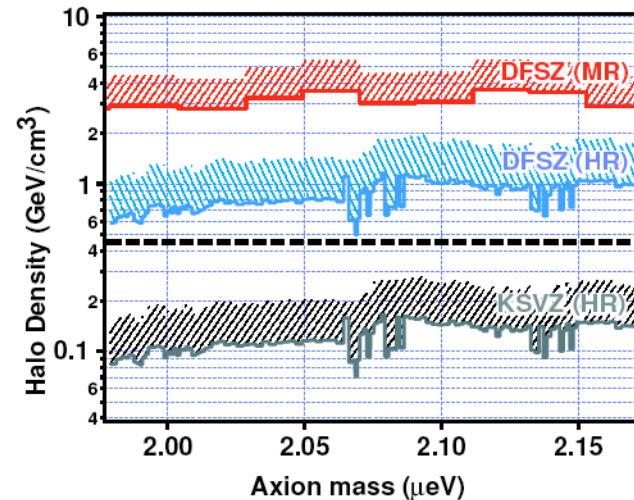
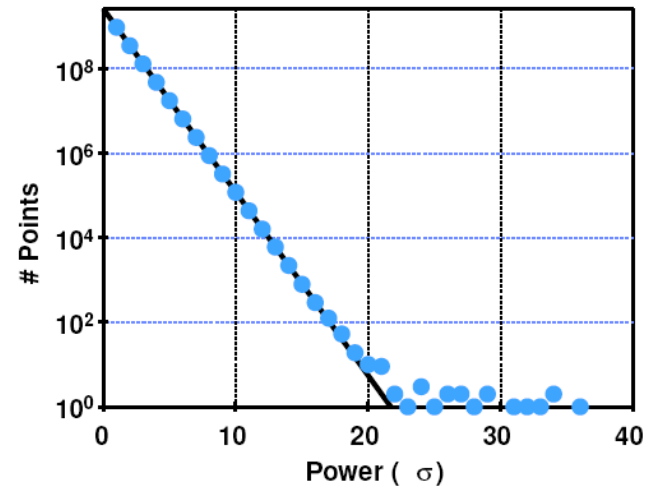
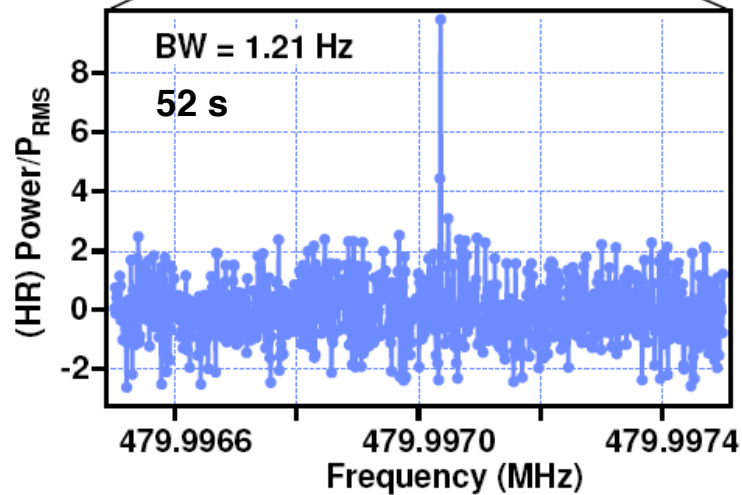
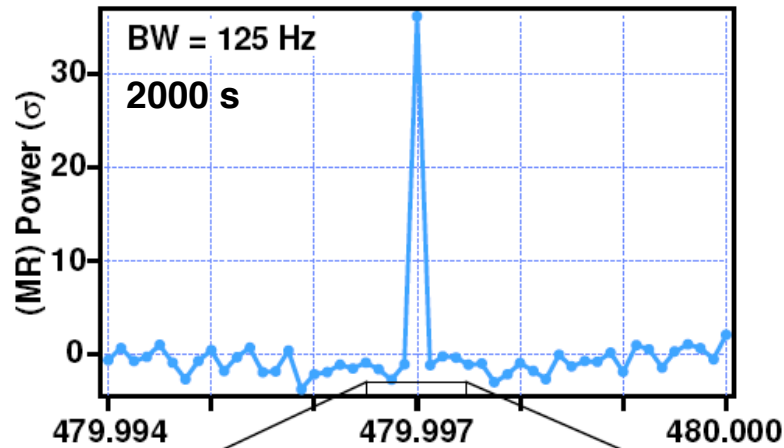


Plausible models have been excluded at the halo density over an octave in mass range

Results of a high-resolution analysis

PRL 95 (9) 091304 (2005)

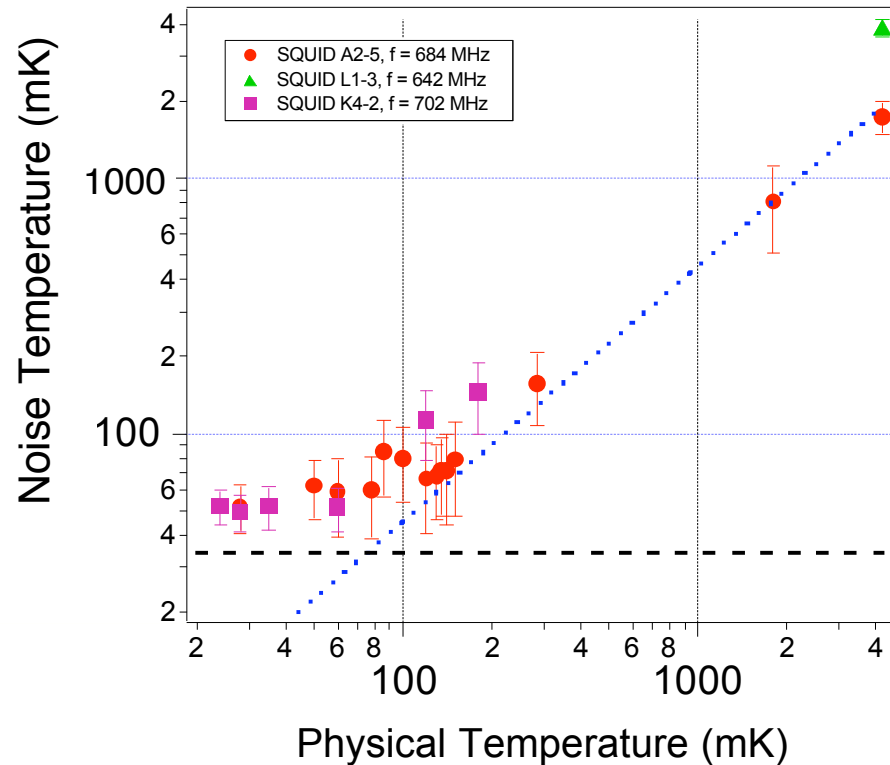
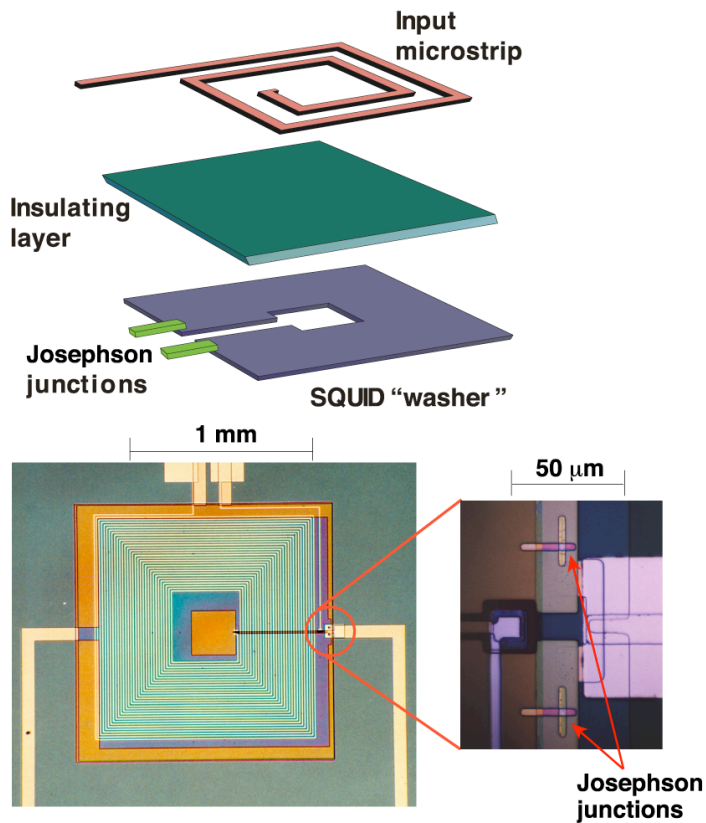
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Measured power in environmental (radio) peak same in Med- & Hi-Res

Upgrade well underway to GHz SQUID amplifiers

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Latest SQUIDs are now within 30% of the Standard Quantum Limit

Rydberg-atom single-quantum detectors

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Atoms with a single electron promoted to a large principal quantum number, $n \gg 1$. Superposition of Rydberg states yields “classical atoms” with macroscopic dimensions (e.g. ~ 1 mm).

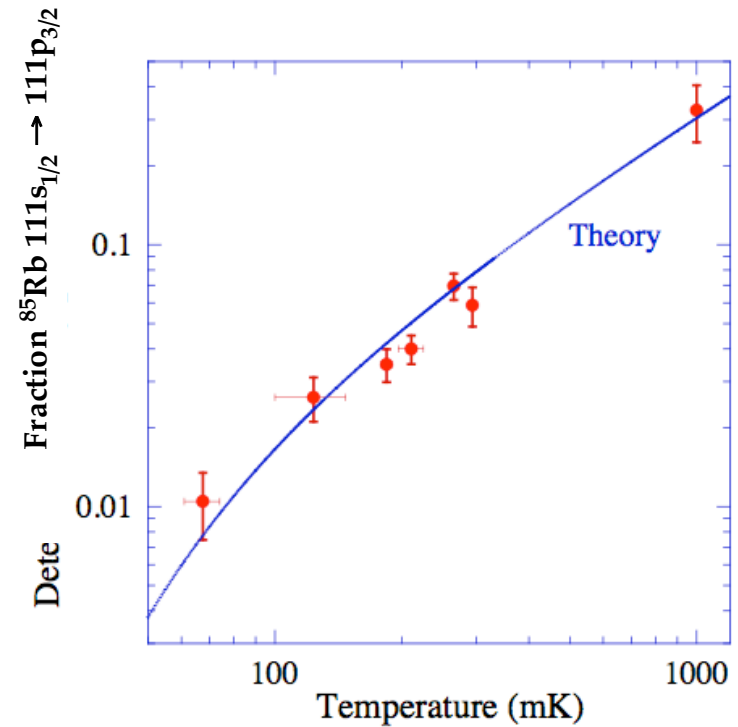
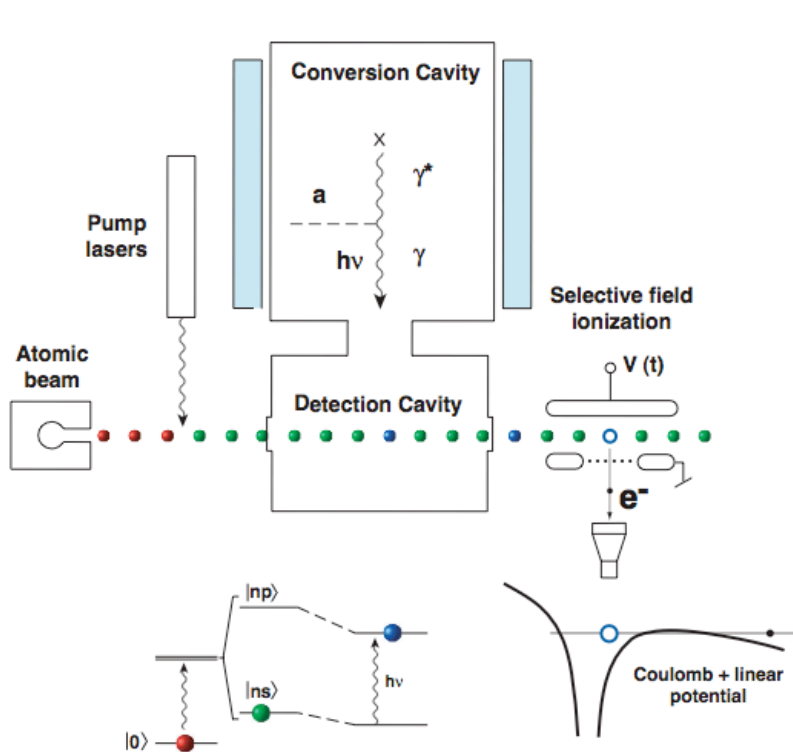
Potential for highly sensitive microwave photon detectors (“RF photo-multiplier tubes”) realized by Kleppner and others in the 1970’s. The axion experiment is an ideal application for Rydberg atoms:

- **Large transition dipole moments** $\langle n \pm 1 | er | n \rangle \propto n^2 a_0$
- **Long lifetimes** $\tau_n \propto n^3$ ($l \ll n$); $\tau_{100} \approx 1$ msec
- **Transitions span microwave range** $\Delta E_n = E_{n+1} - E_n \approx 2R/n^2$; $\Delta E_{100} \approx 7$ GHz

Most importantly, being a phaseless detector (photons-as-particles), the Rydberg-atom detector can evade the standard quantum limit:

$$h\nu = kT$$

Rydberg single-quantum detection (*S. Matsuki et al., Kyoto*)

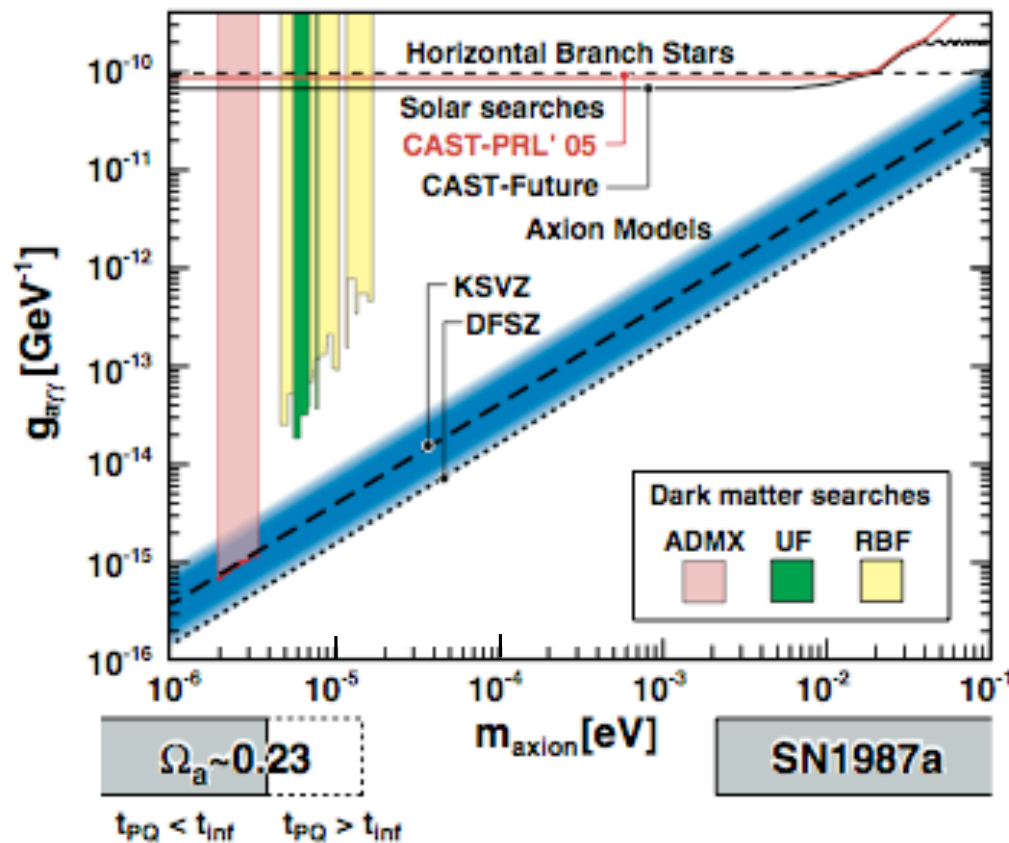


M. Tada et al., Phys. Lett. A (accepted)

The blackbody spectrum has been measured at 2527 MHz a factor of ~ 2 below the standard quantum limit ($\sim 120\text{mK}$)

Summary of axionic dark matter & microwave cavity searches

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Cosmology bounds $m_a > 1 \mu\text{eV}$

- $\Omega_a \propto m_a^{-7/6}$, $\sim O(1)$ for 1-10 μeV
- Why it's a good DM candidate

Astrophysics bounds $m_a < 1 \text{meV}$

- Sn1987a, stellar evolution & lab

Model $g_{a\gamma\gamma}$ banded within ~ 10

- From limited exploration by Kim

ADMX already in region of interest

- SQUIDs enable definitive exp't

The ADMX upgrade is almost complete and will resume operation in 2007

Solar axion searches

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Solar axion spectrum

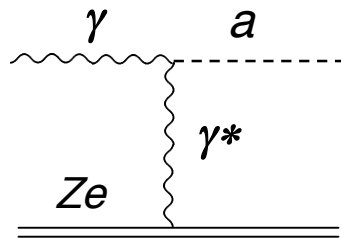
Axion-photon mixing & principle of the experiment

The CERN Axion Search Telescope (CAST)

Results, future plans

The solar axion spectrum

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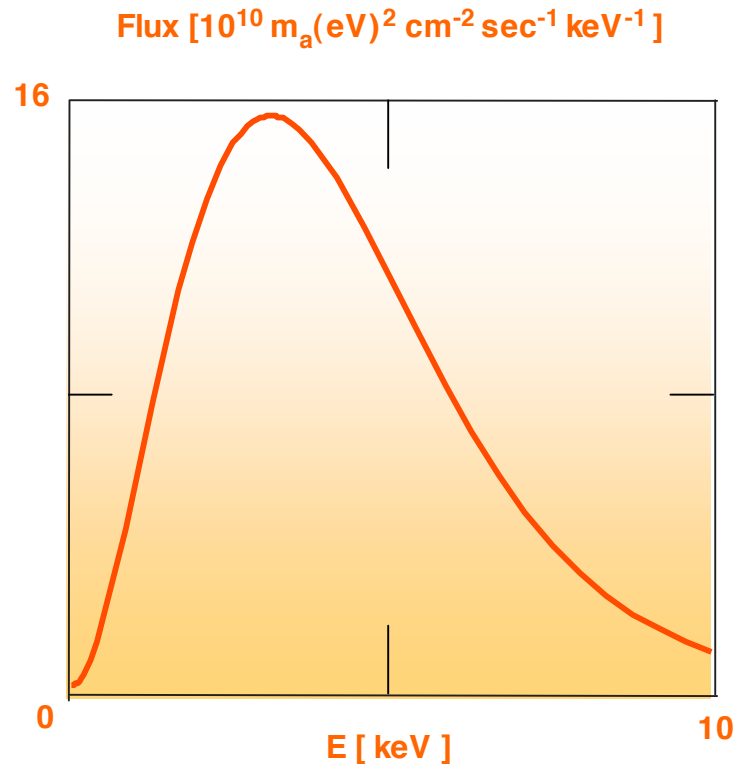
Produced by a Primakoff interaction, with a mean energy of 4.2 keV

$T_{central} = 1.3$ keV, but plasma screening suppresses low energy part of spectrum

The total flux (for KSVZ axions) at the Earth is given by

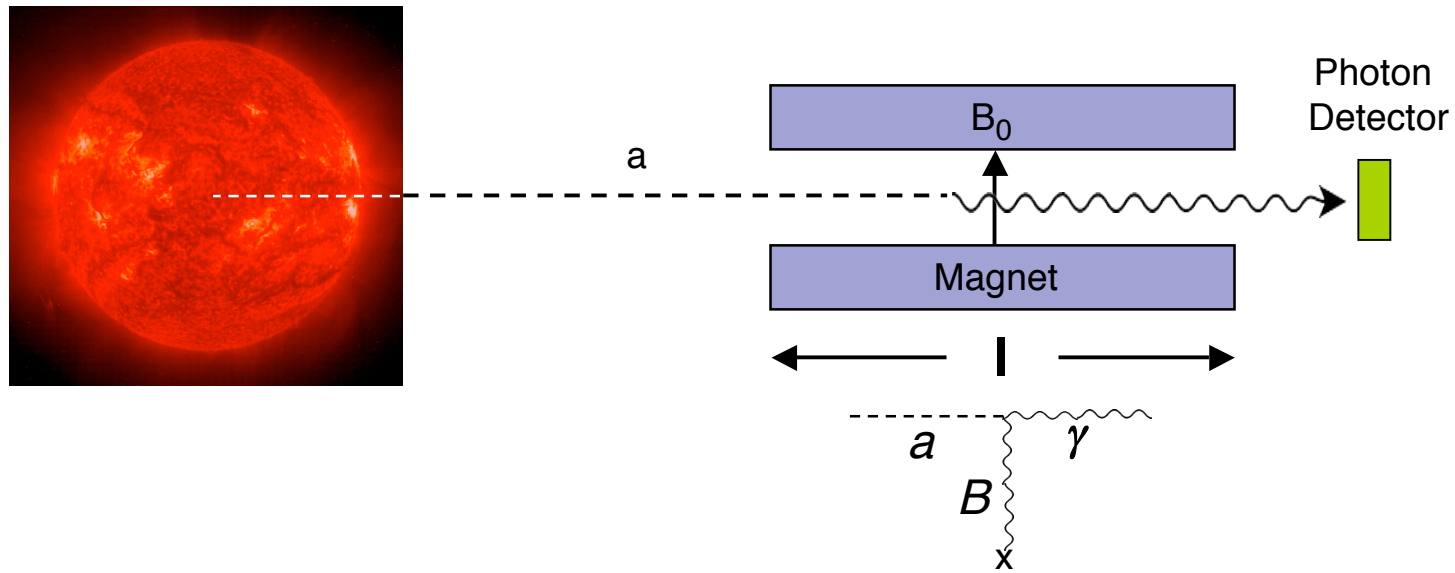
$$\Phi_a = 7.44 \times 10^{11} \text{ cm}^{-2} \text{ sec}^{-1} (m_a / 1\text{eV})^2$$

The dominant contribution is confined to the central 20% of the Sun's radius



Principle of the experiment (*Sikivie's PRL 1983 again!*)

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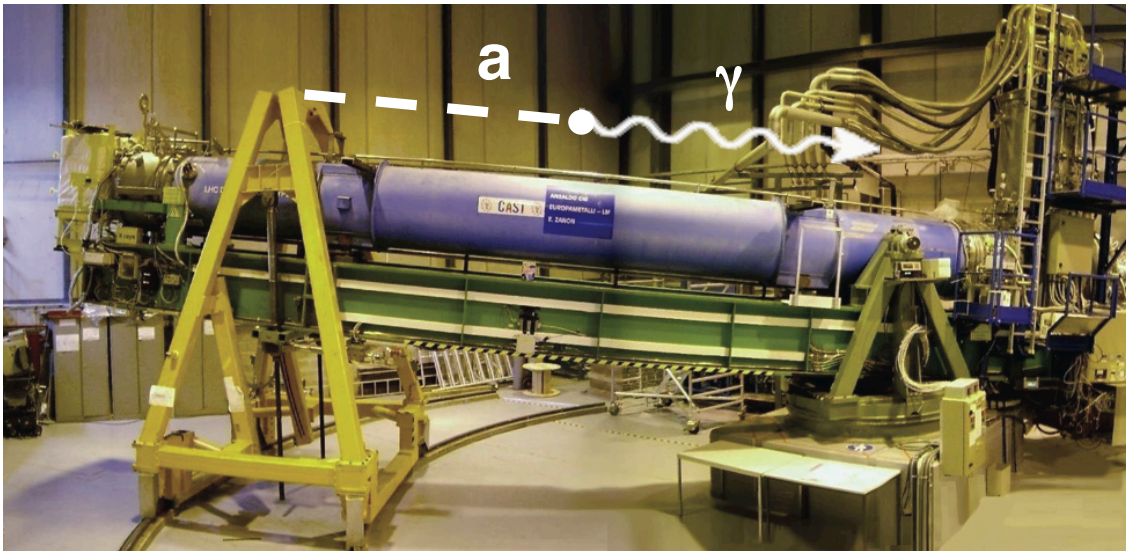


$$L_{a\gamma\gamma} = ag_{a\gamma\gamma} E \cdot B \quad \longrightarrow \quad \Pi(a \leftrightarrow \gamma) = \frac{1}{4} (g_{a\gamma\gamma} B_0 L)^2 |F(q)|^2$$

where $F(q) = \frac{\text{Sin}(qL/2)}{(qL/2)}$, $F(0) = 1$ and $q = k_\gamma - k_a \approx m_a^2 / 2\omega$

The CERN Axion Solar Telescope (CAST)

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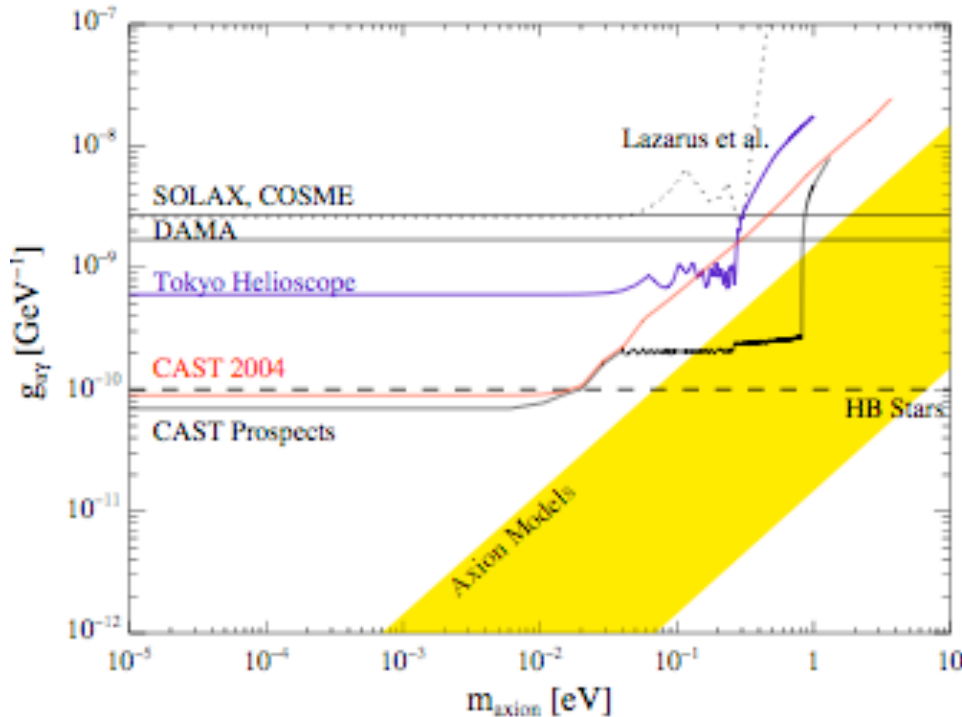


Prototype LHC dipole magnet, double bore, 50 tons, $L \sim 10\text{m}$, $B \sim 10\text{T}$

Tracks the Sun for 1.5 hours at dawn & 1.5 hours at dusk

Instrumented w. 3 technologies: CCD w. x-ray lens; Micromegas; TPC

CAST results and future prospects



CAST has published results equalling the Horizontal Branch Star limit (Red Giant evolution)

The Phase II run underway is pushing the mass limit up into the region of axion models, 0.1-1 eV

Fill the magnet bore with gas (e.g. helium), and tune the pressure

When the plasma frequency equals the axion mass, full coherence and conversion probability are restored:

K. Zioutas et al., Phys. Rev. Lett. **94**, 121301 (2005)

$$\omega_p = (4\pi\alpha N_e / m_e)^{1/2} \equiv m_\gamma$$

KvB et al. PRD 1989

LLNL is providing ^3He for the Phase II run, and fabricating a second x-ray optic

Purely laboratory experiments

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Photon regeneration

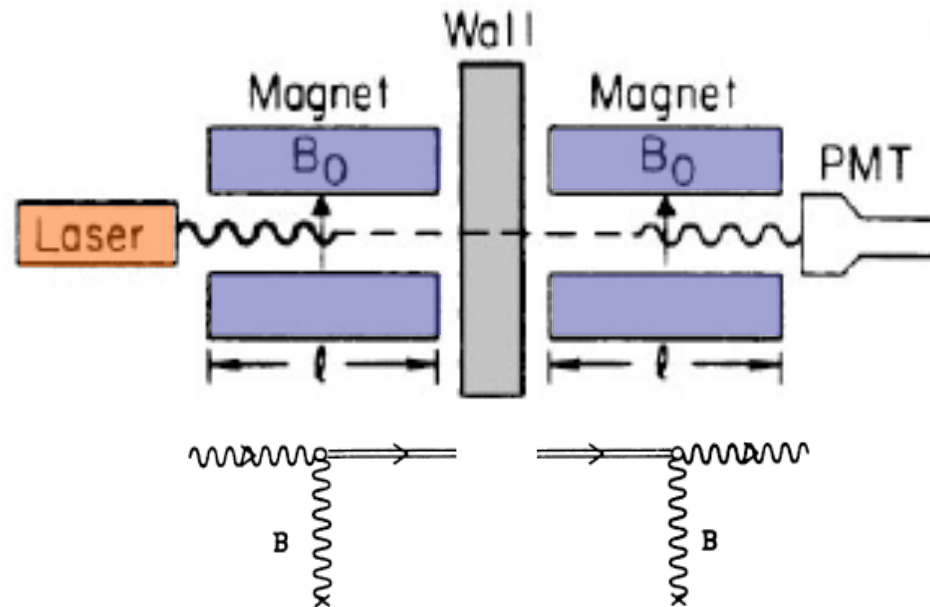
Optical activity of the vacuum

Magnetically-induced vacuum birefringence & dichroism

The PVLAS results

Photon regeneration (a.k.a. “shining light through walls”)

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KvB et al. PRL 59, 759 (1987)

$$P(\gamma \rightarrow a \rightarrow \gamma) = \Pi^2 = 1/16 (gB_0L)^4 |F(q)|^4$$

Difficult to push down to competitive values of the axion-photon coupling

Only measurement to date $g < 7.7 \times 10^{-7} \text{ GeV}^{-1}$ for $m_a < 1 \text{ meV}$ @ BNL

[G. Ruoso et al., Z. Phys. C. 56, 505 (1992)] – but several in preparation now

Vacuum birefringence & dichroism

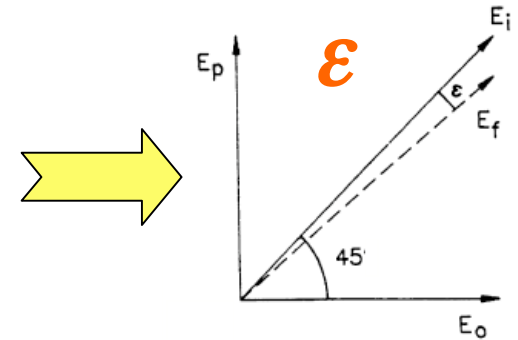
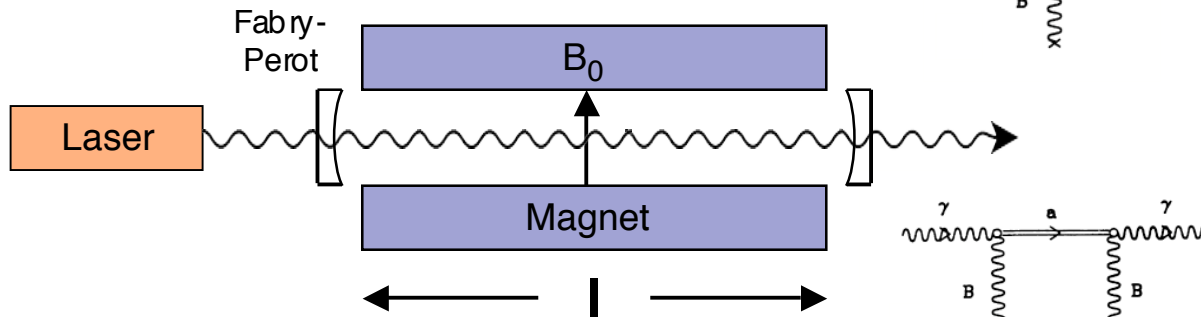
(Maiani, Zavattini, Petronzio, Phys. Lett. 1986)

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Vacuum dichroism

$$e = N \cdot (1/4 g B_0 L)^2 \cdot |F(q)|^2$$

($N = \text{number of passes}$)



Vacuum birefringence

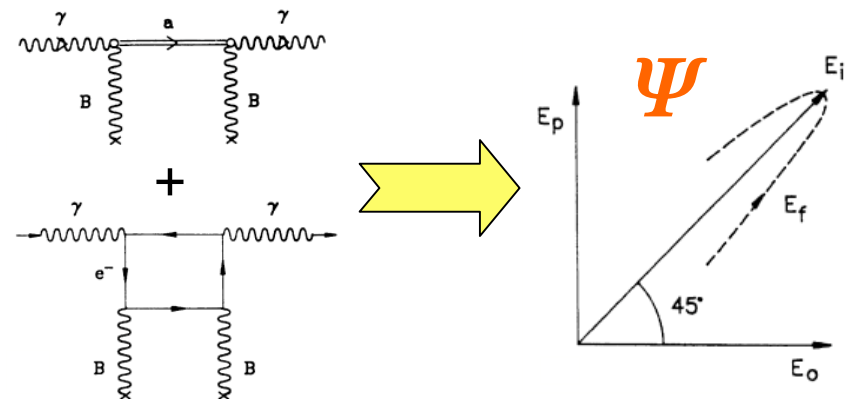
QED:

$$n_{\perp} = 1 + 4/2 \cdot \xi, \quad n_{\parallel} = 1 + 7/2 \cdot \xi$$

$$\xi = \alpha/45\pi (B/B_{\text{crit}})^2, \quad B_{\text{crit}} = m_e^2/e \sim 4.41 \times 10^{13} \text{ G}$$

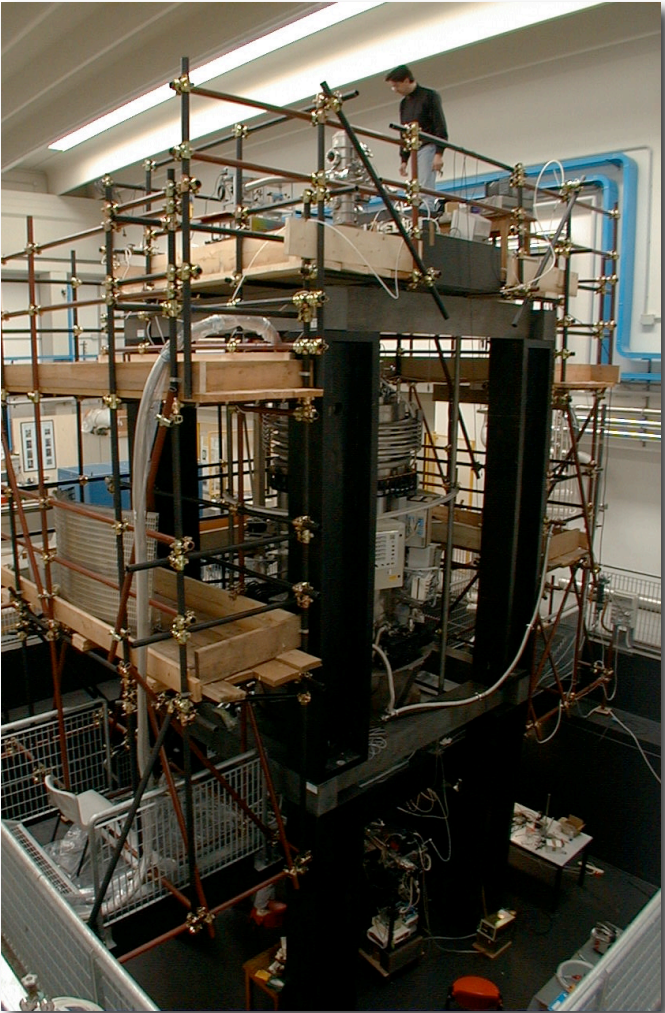
Axion:

$$\Psi = N \cdot (1/96) \cdot (g B_0 m_a)^2 \cdot L^3 / \omega$$

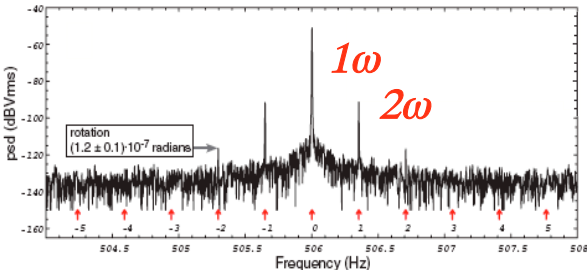
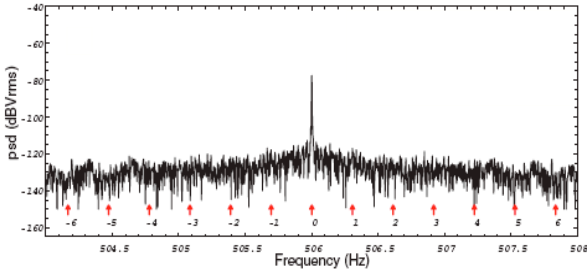
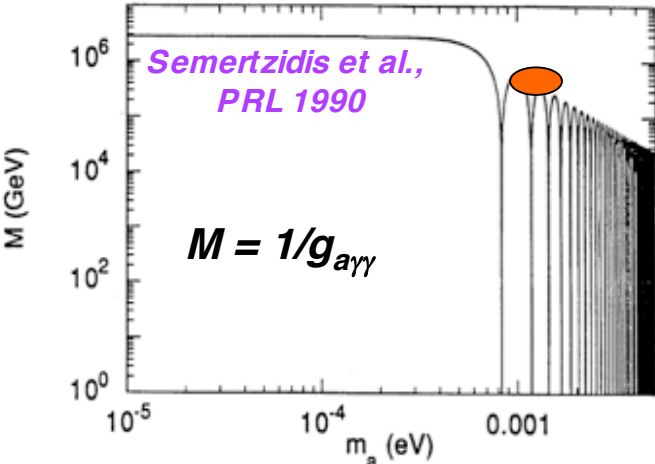


The PVLAS experiment (INFN Legnaro)

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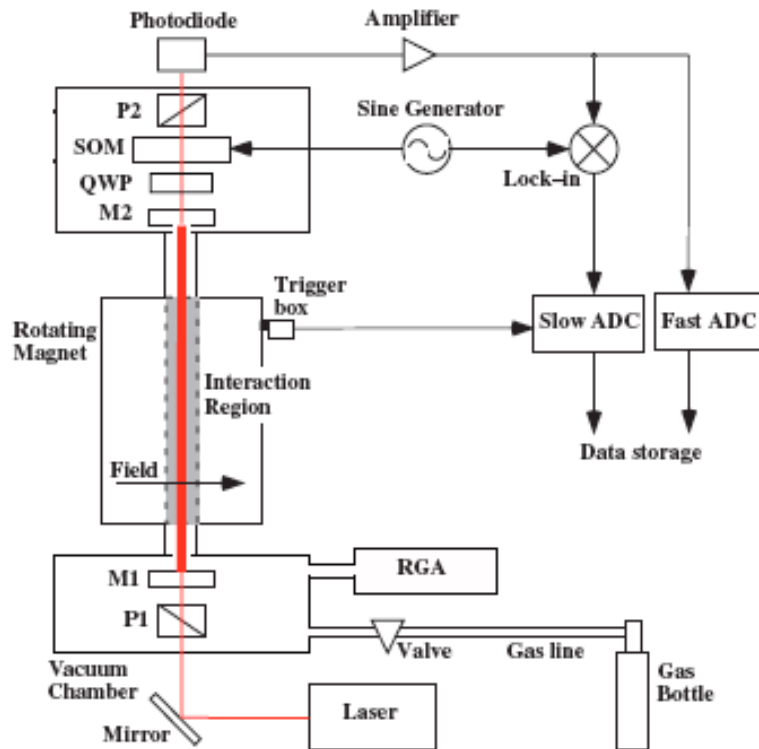


Zavattini et al., PRL 2006

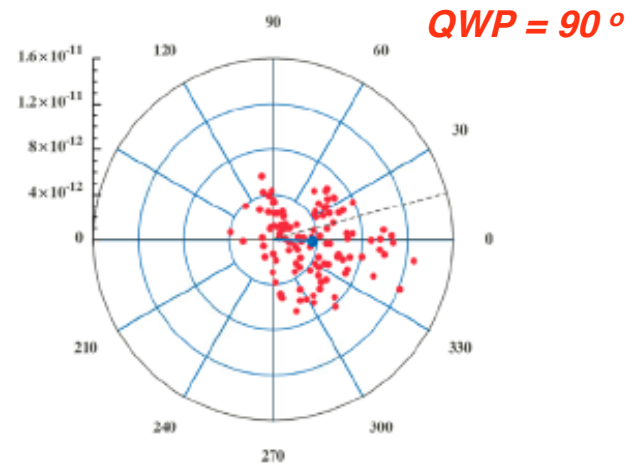
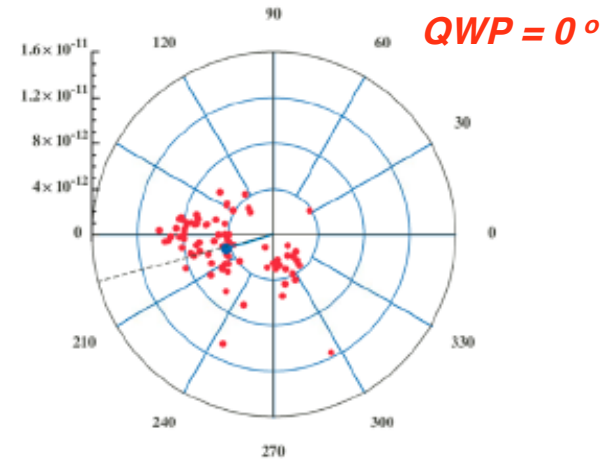


PVLAS details & data

PVLAS Schematic



Phase-Amplitude Plot



The PVLAS results are intriguing but very odd

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The experimenters had hoped to see the QED effect (“light-by-light” scattering), but their sensitivity was not good enough by many orders of magnitude

Their value of $g_{\text{a}\gamma\gamma}$ is ostensibly excluded already by 4 orders of magnitude, by CAST, and stellar evolution (stars would live only a few thousand years)

The allowed region is on the very fringe of the exclusion region of the earlier RBF polarization experiment, plus the photon regeneration experiment

The signal is extremely small: 3.9×10^{-12} rad/pass – the angular width of a pencil lead on the Moon viewed from Earth

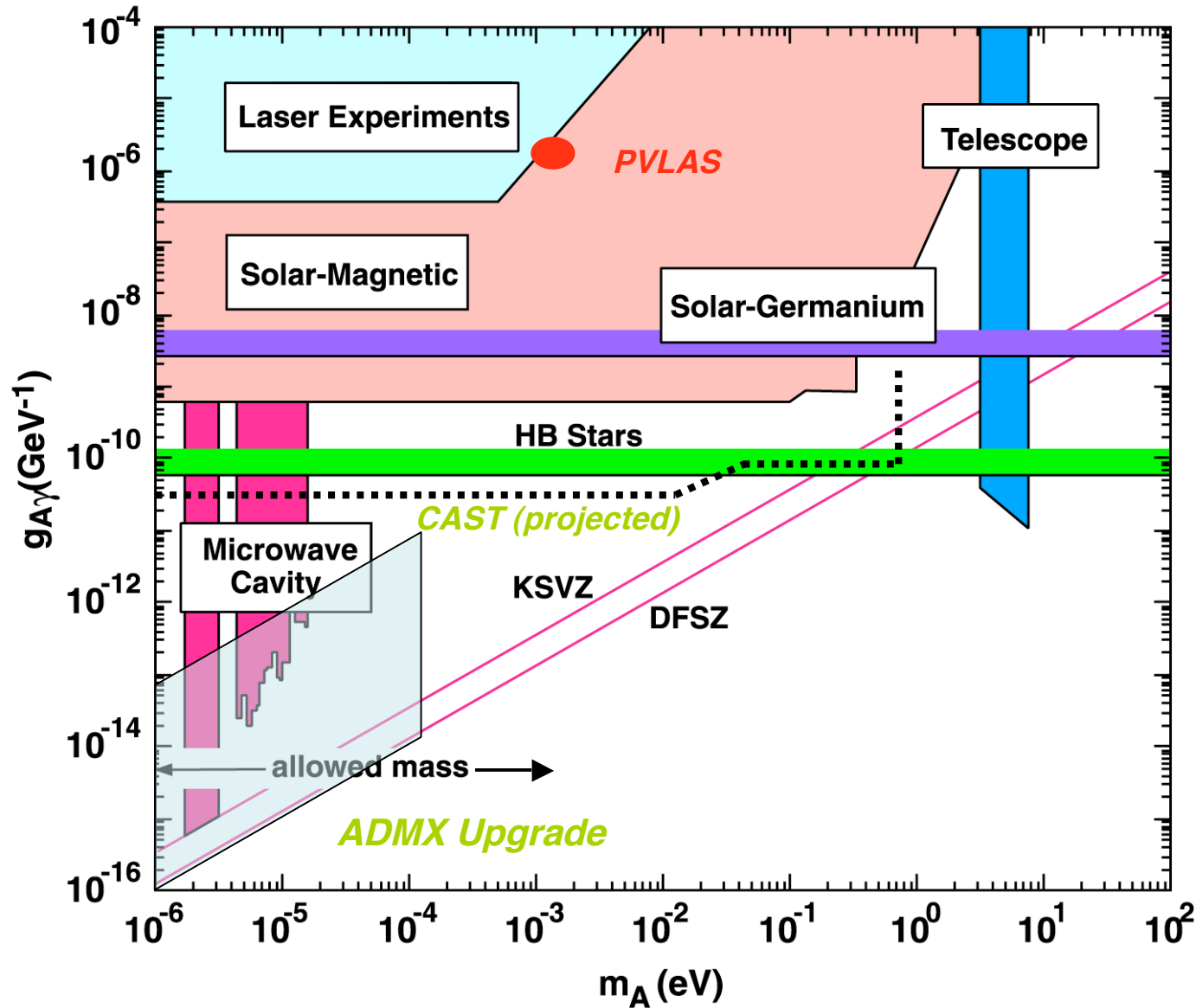
There are evident systematic issues with the experiment: large run-to-run variations in the data, many times the estimated error per point; the unexplained 1ω peak; anomalous dichroism with gases at low pressure, etc.

The effect of stray magnetic fields on the optics, particularly on the Fabry-Perot mirrors may be suspected; this was problematic for the earlier RBF experiment

Nevertheless, this result has launched half a dozen polarization-rotation experiments around the world, and much theoretical work!

Excluded $g_{A\gamma}$ vs. m_A with all experimental and observational constraints

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Summary & final remarks

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The theoretical case is better than ever

- *“If the axion doesn’t exist, please tell me how to solve the Strong-CP problem” (Wilczek)*
- *“Axions may be intrinsic to the structure of string theory” (Witten)*

Experiments are making excellent progress, and discovery would teach us a lot

- *Discovery of dark-matter axions could reveal the detailed history of our galactic evolution*
- *Discovery of solar axions would give us an unprecedented picture of the nuclear-burning core*
- *Discovery of axions in the laboratory would have imponderable consequences*

Experiments have challenges

- *Cavity experiments: With SQUID amps, sensitivity is not an issue, but mass may be*
- *Solar helioscope: The sensitivity will beat the HB limit, but not by much*
- *Lab experiments: While independent of astrophysics/cosmology, the limits are weak*

But remember – Physics is where you find it!

AXION
