

LHC Search of Charged Electroweak Gauge Bosons in Warped Extra Dimensions

GuiYu Huang

University of California, Davis , Department of Physics

LHC LUNCH SEMINAR AT UC DAVIS, PHYSICS

OCTOBER 14, 2008

Kaustubh Agashe, Hooman Davoudiasl, Shrihari Gopalakrishna,
Tao Han, GH, Gilad Perez, Zong-Guo Si, Amarjit Soni
Phys. Rev. D **76**, 115015 (2007) [arXiv:0709.0007](https://arxiv.org/abs/0709.0007)

Agashe, Gopalakrishna, Han, GH, Soni [arXiv:0810.1497](https://arxiv.org/abs/0810.1497)

Kaluza-Klein Theory

pure 5D gravity

$$d\hat{s}^2 = \hat{g}_{\hat{\mu}\hat{\nu}} d\hat{x}^{\hat{\mu}} d\hat{x}^{\hat{\nu}}$$

5 \rightarrow 4 + 1 split, e.g.

$$\hat{g}_{\hat{\mu}\hat{\nu}} = e^{\phi/\sqrt{3}} \begin{pmatrix} g_{\mu\nu} + e^{-\sqrt{3}\phi} A_\mu A_\nu & e^{-\sqrt{3}\phi} A_\mu \\ e^{-\sqrt{3}\phi} A_\nu & e^{-\sqrt{3}\phi} \end{pmatrix}$$

$\hat{\mu} = 0, 1, 2, 3, 4$, and $\mu = 0, 1, 2, 3$

$\hat{x}^{\hat{\mu}} = (x^\mu, y)$

$g_{\mu\nu}(x), A_\mu(x), \phi(x)$

Klein's Compactification

- Non-observation of the extra dimension
- Suppression of field dependence on extra dimension

explained by a compact, small extra dimension

$$M^5 \rightarrow M^4 \times S^1$$

$$g_{\mu\nu}(x, y) = \sum g_{\mu\nu n}(x) e^{inky},$$

$$A_\mu(x, y) = \sum A_{\mu n}(x) e^{inky}$$

$$\phi(x, y) = \sum \phi_n e^{inky}, \quad (kr = 1)$$

\implies Zero modes and infinite towers of excitations

Warped Extra Dimension

Randall-Sundrum Model

- Warping in the 5th Dim:

$$ds^2 = e^{-2k|y|}(\eta_{\mu\nu} dx^\mu dx^\nu) + dy^2$$

- Z_2 Orbifolding: Exponential scale factor between Planck Brane and TeV Brane
- Solves hierarchy problem with a warp factor from AdS background

$$e^{-k\pi R} \sim M_{EW}/M_{pl}$$

$$k\pi R \sim \log M_{pl}/\text{TeV} \sim 34$$

R: size of extra dim.

k: curvature of extra dim.

Warped Extra Dimension

RS in the bulk

- Also explains flavor hierarchy with field profile (or wavefunction) in extra dimension
- Mass and couplings explained by overlappings of profiles

$$C_{mnq}^{FFG} = \int \frac{d\phi}{\sqrt{k}} \frac{e^{t\sigma} \chi_F^{(m)} \chi_F^{(n)} \chi_G^{(q)}}{\sqrt{R}}.$$

Collider Signature of RS

- KK-graviton: dilepton or diphoton signal, need high luminosity
([hep-ph/0006041](#), [hep-ph/0701150](#), [hep-ph/0701186](#))
- KK-gluon: large cross section, $t\bar{t}$ final state
([hep-ph/0612015](#), [hep-ph/0701166](#), [arXiv:0706.3960](#))
- KK-EW gauge boson: dominated by $t\bar{t}$. $\ell^+\ell^-$ suppressed. Other channels?
- KK-fermion ([arXiv:0706.1281](#) [arXiv:0712.0095](#))

Gauge Symmetry Breaking

Bulk Gauge Group

Neglecting $SU(3)_C$,

$$\begin{aligned}
 & SU(2)_L \times SU(2)_R \times U(1)_X \\
 \Rightarrow & \text{(viaBoundaryCondition)} \quad SU(2)_L \times U(1)_Y \\
 \Rightarrow & \text{(viaHiggsMechanism)} \quad U(1)_Q
 \end{aligned}$$

$$Y = T_{3R} + X, \quad Q = T_{3L} + Y/2.$$

Gauge Bosons

- Charged: W_L^\pm, W_R^\pm

- Neutral: W_L^3, W_R^3, X

$$\Rightarrow W_L^3, B, Z_X$$

$$\Rightarrow A, Z, Z_X$$

Only A, Z, W_L have zero-mode, W_R, Z_X don't. (--) BC

Representations

Fermions

- $Q_L = (2, 2) = \begin{pmatrix} t_L & \chi_L \\ b_L & T_L \end{pmatrix}$
- $t_R = (1, 1) \text{ or } (1, 3) \in (1, 3) \oplus (3, 1) = \begin{pmatrix} \chi_R'' \\ t_R \\ B_R'' \end{pmatrix} \oplus \begin{pmatrix} \chi_R''' \\ T_R''' \\ B_R''' \end{pmatrix}$

Higgs

- $\Sigma = (2, 2)$

For these reps

- $Zb\bar{b}$: protected by custodial symmetry
[Agashe, Contino, DaRold, Pomarol - 06]
- Precision EW constraints $\Rightarrow M_{Z'} \gtrsim 2 - 3\text{TeV}$
[Carena, Ponton, Santiago, Wagner - 06,07]

Two Cases We Consider

- Case (i): $t_R \rightarrow (1, 1)$, $c_{Q_L^3} = 0$, $c_{t_R} = 0.4$.
- Case (ii): $t_R \rightarrow (1, 3)$, $c_{Q_L^3} = 0.4$, $c_{t_R} = 0$.

Charged Gauge Bosons

Mass Spectrum

Mass term for zero mode and 1st excitations:

$$\begin{pmatrix} W_L^{+(0)} & W_{L1}^+ & W_{R1}^+ \end{pmatrix} \mathcal{M}_W^2 \begin{pmatrix} W_L^{-(0)} \\ W_{L1}^- \\ W_{R1}^- \end{pmatrix}$$

$$\mathcal{M}_W^2 =$$

$$\begin{pmatrix} m_W^2 & m_W^2 \sqrt{k\pi R} & -m_W^2 \sqrt{k\pi R} \frac{g_R}{g} \\ m_W^2 \sqrt{k\pi R} & m_{KK}^2 + m_W^2 k\pi R & -m_W^2 k\pi R \frac{g_R}{g} \\ -m_W^2 \sqrt{k\pi R} \frac{g_R}{g_L} & -m_W^2 k\pi R \frac{g_R}{g} & 0.963 m_{KK}^2 + m_W^2 k\pi R \left(\frac{g_R}{g}\right)^2 \end{pmatrix}$$

$$m_{W_{1L}} \sim m_{KK}, \quad m_{W_{1R}} \sim 0.981 m_{KK}$$

Mixings

$$g' = \frac{g_X g_R}{\sqrt{g_R^2 + g_X^2}}, \quad s' = \frac{g_X}{\sqrt{g_R^2 + g_X^2}}, \quad c' = \sqrt{1 - s'^2},$$

$$e = \frac{g_L g'}{\sqrt{g'^2 + g_L^2}}, \quad s_W = \frac{g'}{\sqrt{g'^2 + g_L^2}}, \quad c_W = \sqrt{1 - s_W^2},$$

$$g_Z = g_L / c_W, \quad g_{Z'} = g_R / c'.$$

For $g_R = g_L$, $s' = 0.55$, $c' = 0.84$.

$$\sin \theta_{01} \approx \left(\frac{M_Z}{M_{Z_1}} \right)^2 \sqrt{k\pi R},$$

$$\sin \theta_{01X} \approx - \left(\frac{M_Z}{M_{Z_{X1}}} \right)^2 \left(\frac{g_{Z'}}{g_Z} \right) c'^2 \sqrt{k\pi R}.$$

$$\sin \theta_{0L} \approx \left(\frac{M_W}{M_{W_{L1}}} \right)^2 \sqrt{k\pi r_c},$$

$$\sin \theta_{0R} \approx - \left(\frac{M_W}{M_{W_{R1}}} \right)^2 \left(\frac{g_R}{g_L} \right) \sqrt{k\pi r_c}.$$

For $m_{KK} = 2 \text{ TeV}$

$$s_{01} \sim 0.013, \quad s_{01X} \sim -0.01, \quad s_{0L} \sim 0.01, \quad s_{0R} \sim -0.01.$$

Mixings

$$\tan 2\theta_1 = \frac{-2M_Z^2(g_{Z'}/g_Z)c'^2 k_{\pi R}}{(M_{Z_{X1}}^2 - M_{Z_1}^2) + M_Z^2((g_{Z'}/g_Z)^2 c'^4 - 1) k_{\pi R}}.$$

$$\tan 2\theta_1^c = \frac{-2M_W^2(g_R/g_L)k_{\pi R}}{(M_{W_{R1}}^2 - M_{W_{L1}}^2) + M_W^2((g_R/g_L)^2 - 1) k_{\pi R}}.$$

For $m_{KK} = 2000$ GeV

$$s_1 = 0.48, c_1 = 0.88; \quad s_1^c = 0.6, c_1^c = 0.8.$$

W' Couplings

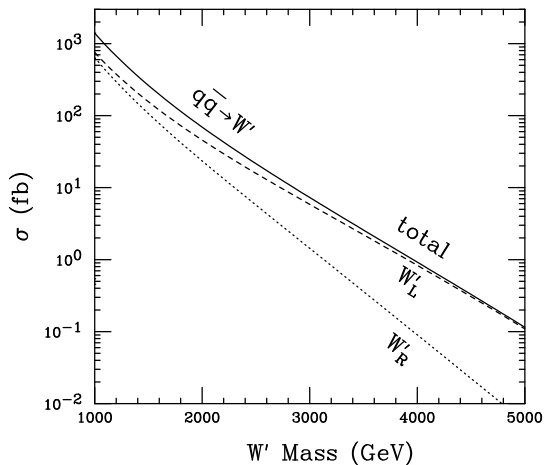
$$\begin{aligned}
 \frac{g_{RS}^{q\bar{q}, \ell\bar{\ell}} A^{(1)}}{g_{SM}} &\simeq -\xi^{-1} \approx -\frac{1}{5} \\
 \frac{g_{RS}^{Q^3 \bar{Q}^3 A^{(1)}}}{g_{SM}}, \frac{g_{RS}^{t_R \bar{t}_R A^{(1)}}}{g_{SM}} &\simeq 1 \text{ to } \xi (\approx 5) \\
 \frac{g_{RS}^{HHA^{(1)}}}{g_{SM}} &\simeq \xi \approx 5 \quad (H = h, W_L, Z_L) \\
 \frac{g_{RS}^{A^{(0)} A^{(0)} A^{(1)}}}{g_{SM}} &\sim 0
 \end{aligned}$$

$c_{Q_L^3} = 0, c_{t_R} = 0.4$	Q_L^3	t_R	other fermions
$\mathcal{I}_{++,++}^{++}$	$-\frac{1.13}{\xi} + 0.7\xi \approx 3.9$	$-\frac{1.13}{\xi} + 0.2\xi \approx 1$	$-\frac{1.13}{\xi} \approx -0.2$
$\mathcal{I}_{-+,-+}^{++}$	ξ	ξ	$-\frac{1.13}{\xi} \approx -0.2$
$\mathcal{I}_{++,,-+}^{-+}$	$0.8\xi \approx 4.6$	$0.4\xi \approx 2.3$	≈ 0
$c_{Q_L^3} = 0.4, c_{t_R} = 0$	Q_L^3	t_R	other fermions
$\mathcal{I}_{++,++}^{++}$	$-\frac{1.13}{\xi} + 0.2\xi \approx 1$	$-\frac{1.13}{\xi} + 0.7\xi \approx 3.9$	$-\frac{1.13}{\xi} \approx -0.2$
$\mathcal{I}_{-+,-+}^{++}$	ξ	ξ	$-\frac{1.13}{\xi} \approx -0.2$
$\mathcal{I}_{++,,-+}^{-+}$	$0.4\xi \approx 2.3$	$0.8\xi \approx 4.6$	≈ 0

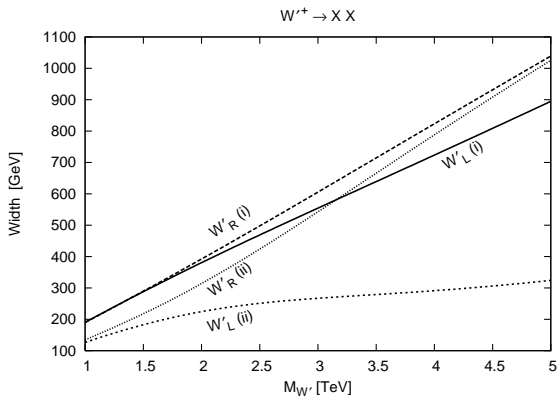
Table: Values of $\psi\psi W'^{\pm}$ overlap integrals. $\xi = \sqrt{k\pi r_c} = 5.83$. All SM fermions have (++) BC, "exotic" BSM fermions have (-+), W_{L1} has (++) BC, and, W_{R1} has (-+) BC.

W' Drell-Yan Production

Total W' Cross Section at LHC



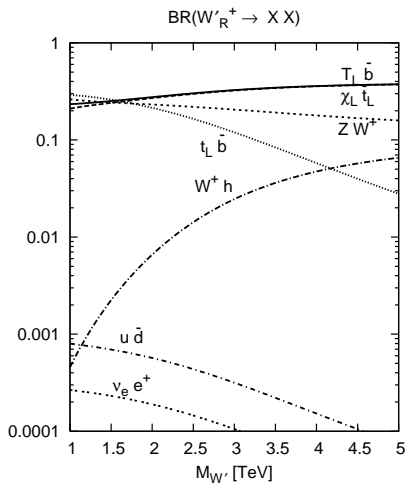
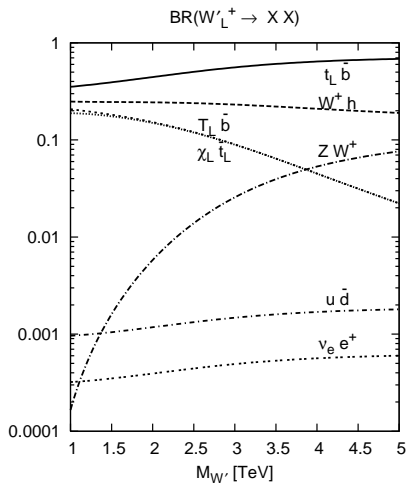
W' Decay Widths



Case (i): $t_R \rightarrow (1, 1)$, $c_{Q_L^3} = 0$ and $c_{t_R} = 0.4$. All the other c 's > 0.5

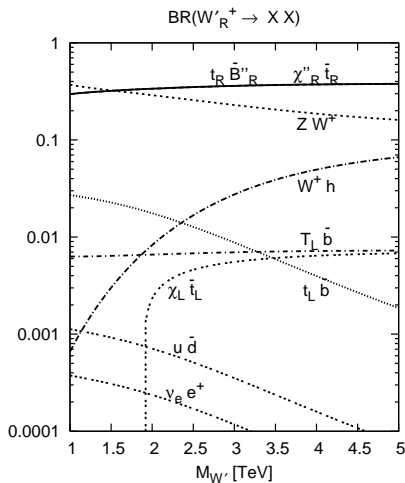
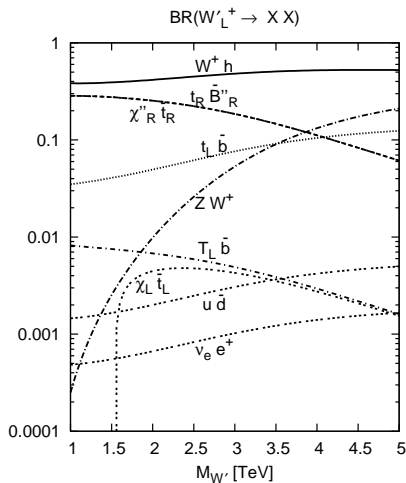
Case (ii): $t_R \rightarrow (1, 3)$, $c_{Q_L^3} = 0.4$ and $c_{t_R} = 0$. All the other c 's > 0.5

W' Branchings



Case (i)

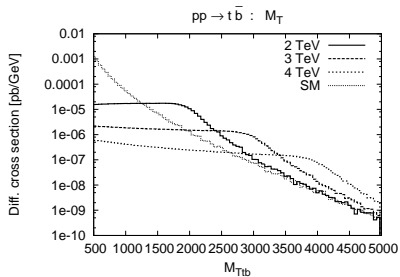
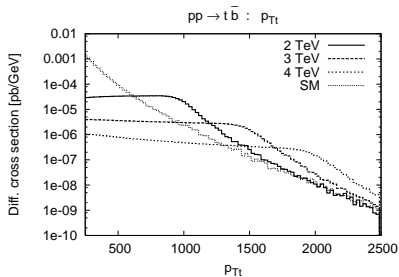
W' Branchings



Case (ii)

$$W'^+ \rightarrow t\bar{b}$$

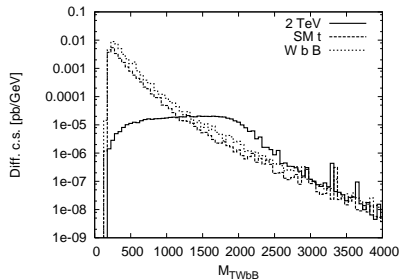
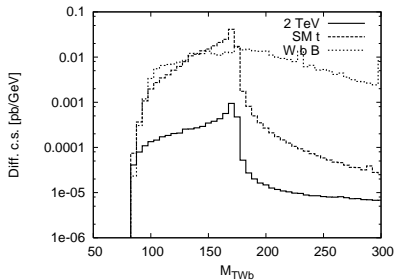
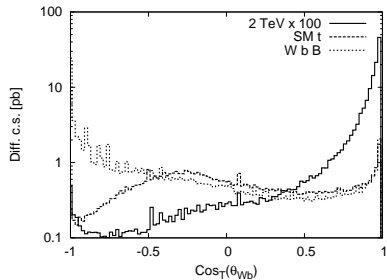
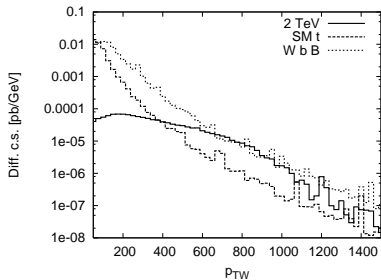
$$pp \rightarrow W'^+ \rightarrow t\bar{b} \rightarrow b\bar{b}\ell\bar{\nu}$$



$$M_{TWb} = \left(\sqrt{p_{TW}^2 + m_W^2} + p_{Tb} \right)^2 - |\mathbf{p}_{TW} + \mathbf{p}_{Tb}|^2,$$

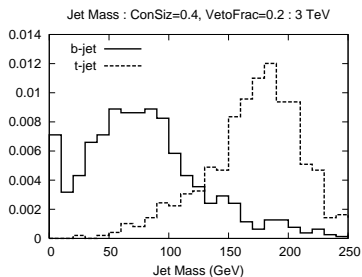
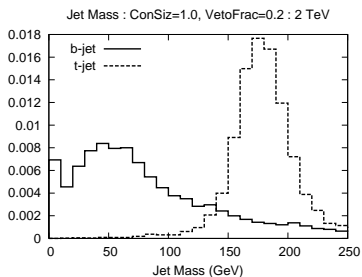
$$M_{TWb\bar{b}} = p_{Tb} + p_{T\bar{b}} + \sqrt{p_{TW}^2 + m_W^2}.$$

$$W'^+ \rightarrow t\bar{b}$$



$$W'^+ \rightarrow t\bar{b}$$

To contend with $t\bar{t}$ background from SM and KK gluons:



Jet masses of top and bottom jets.

$$W'^+ \rightarrow t\bar{b}$$

2 TeV	Basic	Wb cuts	b-tag	M_{TWbb}	j_M	# Evt	S/B	S/ \sqrt{B}
Sig (i)	8.9	7	1.1	0.44	0.2	20	2.5 (1.4)	7 (5.3)
SM top	1431	372	60	0.09	0.04	4		
SM Wbb	517	66	10.6	9×10^{-3}	4×10^{-3}	0.4		
SM Wbj	9×10^3	2×10^3	20	0.04	0.02	2		
SM Wcj	4×10^3	700	4	10^{-3}	0.5×10^{-3}	0.05		
SM Wjj	2×10^5	2×10^4	13	0.03	0.01	1		
SM $t\bar{t}$	4×10^4	10^4	2×10^3	4.5	0.02	2		
$G^{(1)} t\bar{t}$ (i)	246	188	30	10	0.04	4		

3 TeV	Basic	Wb cuts	b-tag	M_{TWbb}	j_M	# Evt	S/B	CL
Sig (i)	1.5	1.1	0.18	0.04	0.02	7	5.8 (0.9)	0.995 (0.95)
SM top	1431	372	60	4×10^{-3}	2×10^{-3}	0.6		
SM Wbb	517	66	10.6	4×10^{-4}	2.3×10^{-4}	0.07		
SM Wbj	9×10^3	2×10^3	20	10^{-3}	0.5×10^{-3}	0.2		
SM Wcj	4×10^3	700	4	10^{-4}	0.5×10^{-4}	0.02		
SM Wjj	2×10^5	2×10^4	13	2×10^{-3}	10^{-3}	0.3		
SM $t\bar{t}$	4×10^4	10^4	2×10^3	0.21	5.3×10^{-3}	1.6		
$G^{(1)} t\bar{t}$ (i)	32	24	4	0.64	0.02	5		

Case (i). $\mathcal{L} = 100 \text{ fb}^{-1}$ for 2 TeV and $\mathcal{L} = 300 \text{ fb}^{-1}$ for 3 TeV

$$W'^+ \rightarrow t\bar{b}$$

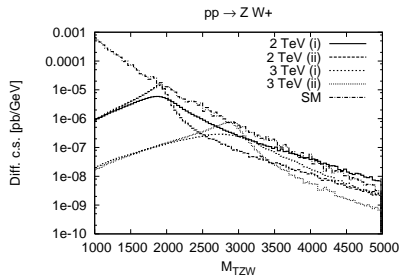
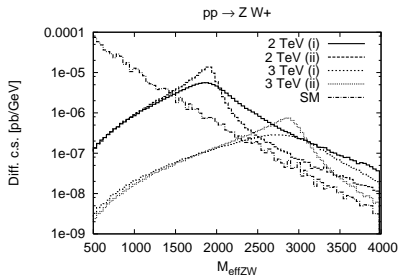
2 TeV	Basic	Wb cuts	b-tag	M_{TWbb}	j_M	# Evt	S/B	S/\sqrt{B}
Sig (ii)	0.75	0.6	0.1	0.05	0.03	30	0.38 (0.2)	3.4 (2.5)
SM top	1431	372	60	0.09	0.04	40		
SM $W_{j_1 j_2}$	2.1×10^5	2.2×10^4	48	0.08	0.04	40		
SM $t\bar{t}$	4×10^4	10^4	2×10^3	4.5	0.02	20		
$G^{(1)} t\bar{t}$ (ii)	207	180	29	12.8	0.05	50		

Case (ii). $\mathcal{L} = 1000 \text{ fb}^{-1}$ for 2 TeV
 Small coupling to top

$W' \rightarrow ZW$

- fully leptonic. $Z \rightarrow \ell\ell$, $W \rightarrow \ell\nu$
- semi-leptonic.
 - $Z \rightarrow \ell\ell$, $W \rightarrow jj$
 - $Z \rightarrow jj$, $W \rightarrow \ell\nu$
- fully hadronic. $Z \rightarrow jj$, $W \rightarrow jj$. (BG too large)

W' → ZW



$$M_{\text{eff}ZW} = p_{TZ} + p_{TW},$$

$$M_{TZW} = \sqrt{p_{TZ}^2 + M_Z^2} + \sqrt{p_{TW}^2 + M_W^2}.$$

$$W' \rightarrow ZW \rightarrow (\ell\ell)(\ell\nu)$$

2 TeV	Basic	M_{eff}	M_T	\mathcal{L}	# Evts	S/B	CL
Signal (i)	0.13	0.13	0.1	100	10	5	0.9995
Signal (ii)	0.17	0.16	0.13	100	13	6.5	> 0.9995
SM Z W	42	0.16	0.02		2		

3 TeV	Basic	M_{eff}	M_T	\mathcal{L}	# Evts	S/B	CL
Signal (i)	0.01	0.01	0.006	1000	6	6	0.99
Signal (ii)	0.014	0.01	0.01	1000	10	10	> 0.9995
SM Z W	42	0.05	0.001		1		

$M_{eff} > 1$ TeV (for 2 TeV) and $M_{eff} > 1.25$ TeV (for 3 TeV).

$1.5 < M_{TZW} < 2.5$ TeV (for 2 TeV) and

$2.4 < M_{TZW} < 3.6$ TeV (for 3 TeV).

Clean signal, large S/B, but limited statistics.

$$W' \rightarrow ZW \rightarrow (\ell\ell)(jj)$$

2 TeV	Basic	M_{eff}	M_{inv}	M_{jet}	\mathcal{L}	# Evts	S/B	S/\sqrt{B}
Sig (i)	0.4	0.4	0.16	0.13	1000	130	0.2	5
Sig (ii)	0.5	0.48	0.38	0.3	300	90	0.5	6.4
SM ZW	128	0.5	0.05	0.04		40, 12		
SM Z + 1j	3580	63	2.1	0.63		630, 189		

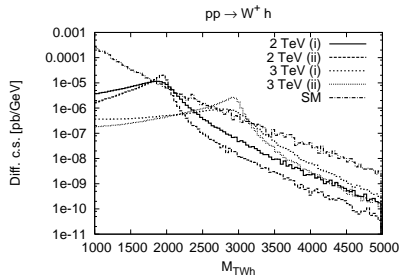
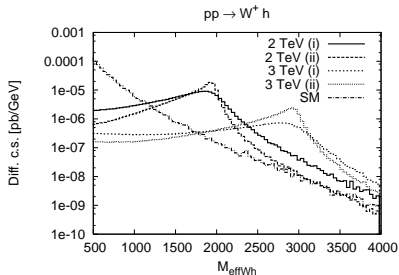
3 TeV	Basic	M_{eff}	M_{inv}	M_{jet}	\mathcal{L}	# Evts	S/B	S/\sqrt{B}
Sig (i)	0.03	0.03	0.01	—	1000	10	0.07	0.8
Sig (ii)	0.04	0.04	0.03	—	1000	30	0.22	2.6
SM ZW	128	0.16	0.006	—		6		
SM Z + 1j	3580	25	0.13	—		130		

$$W' \rightarrow ZW \rightarrow (jj)(\ell\nu)$$

2 TeV	Basic	M_{eff}	M_T	M_{jet}	\mathcal{L}	# Evts	S/B	S/\sqrt{B}
Sig (i)	1	1	0.38	0.3	1000	300	0.1	5.3
Sig (ii)	1.3	1.2	0.64	0.5	300	150	0.16	4.9
SM ZW	318	1.2	0.04	0.03		30, 9		
SM W + 1j	3.1×10^4	224	10.5	3.15		3150, 945		

3 TeV	Basic	M_{eff}	M_T	M_{jet}	\mathcal{L}	# Evts	S/B	S/\sqrt{B}
Sig (i)	0.08	0.08	0.016	–	1000	16	0.02	0.6
Sig (ii)	0.1	0.1	0.04	–	1000	40	0.06	1.5
SM ZW	318	0.4	0.002	–		2		
SM W + 1j	3.1×10^4	89	0.68	–		680		

$W' \rightarrow Wh$



$$M_{\text{eff}Wh} = p_{TW} + p_{Th},$$

$$M_{TW^+h} = \sqrt{p_{TW}^2 + M_W^2} + \sqrt{p_{Th}^2 + m_h^2}.$$

$m_h = 120 \text{ GeV}: h \rightarrow b\bar{b}, W \rightarrow \ell\nu$

$m_h = 150 \text{ GeV}: h \rightarrow WW \rightarrow \ell\nu jj, W \rightarrow jj$

$$W' \rightarrow Wh \rightarrow \ell E_T b \bar{b}$$

2 TeV	Basic	M_{eff}	M_T	b-tag	\mathcal{L}	# Evts	S/B	S/\sqrt{B}
Sig (i)	1.8	1.46	0.55	0.35	100	35	0.65	4.8
Sig (ii)	1.64	1.5	0.8	0.5	100	50	1	6.4
SM Wh	42.9	0.35	0.016	0.01		1		
SM $W + 1j$	3.1×10^4	224	10.5	0.53		53		

3 TeV	Basic	M_{eff}	M_T	b-tag	\mathcal{L}	# Evts	S/B	CL
Sig (i)	0.26	0.19	0.04	0.03	300	9	1	0.99
Sig (ii)	0.33	0.3	0.12	0.08	300	24	2.4	> 0.9995
SM Wh	42.9	0.13	0.001	6×10^{-4}		0.2		
SM $W + 1j$	3.1×10^4	89	0.68	0.03		9		

$$W' \rightarrow Wh \rightarrow (jj)WW \rightarrow (jj)\ell E_T(jj)$$

2 TeV	Basic	M_{eff}	M_T	M_{jet}	\mathcal{L}	# Evts	S/B	CL
Sig (i)	1.57	1.27	0.43	0.34	100	34	4	$\gg 0.9995$
Sig (ii)	2.1	1.9	0.9	0.7	100	70	7	$\gg 0.9995$
SM Wh	25.63	0.31	0.014	0.01		1		
SM $h+1j$	222.6	1.97	0.07	0.02		2		
SM $W+2j$	3×10^4	35.5	0.62	0.06		6		

3 TeV	Basic	M_{eff}	M_T	M_{jet}	\mathcal{L}	# Evts	S/B	CL
Sig (i)	0.22	0.17	0.04	0.035	300	11	2	0.9987
Sig (ii)	0.3	0.26	0.1	0.09	300	27	4	$\gg 0.9995$
SM Wh	25.63	0.12	8×10^{-4}	7×10^{-4}		0.2		
SM $h+1j$	222.6	0.72	5×10^{-3}	2×10^{-3}		0.6		
SM $W+2j$	3×10^4	4.1	0.05	0.015		4.5		

$$W' \rightarrow W \rightarrow \ell E_T$$

2 TeV	Basic	M_{eff}	M_T
Sig (i)	0.04	0.024	0.012
Sig (ii)	0.05	0.04	0.02
SM W	4×10^3	6.9	0.44

Branching too small.

RS W' at LHC

- Interesting phenomenology in multiple channels
- LHC a discovery machine for RS weak bosons of mass $\sim 2 - 3 \text{ TeV}$ with integrated luminosity $\sim 100 - 1000 \text{ fb}^{-1}$
- Higher mass range may require upgrade in luminosity
- Though some channels may be explored simultaneously, hard to extract precise coupling information
- Comparison to RS Z'
- Comparisons to 4D dual models?