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QM 115B Loc. 22
only nonzero matrix elements

$$\langle 100 | z | 210 \rangle \quad \langle 100 | x | 21\pm 1 \rangle \quad \langle 100 | y | 21\pm 1 \rangle$$

$$\langle 100 | z | 210 \rangle = \left(\frac{2^7 \sqrt{2}}{3^5} \right) a$$

$$\langle 100 | x | 21\pm 1 \rangle = \mp \frac{2^7}{3^5} a$$

$$\langle 100 | y | 21\pm 1 \rangle = -i \frac{2^7}{3^5} a$$

$$\langle 100 | \vec{r} | 210 \rangle = 0 \quad \langle 100 | \vec{r} | 210 \rangle = \frac{2^7 \sqrt{2}}{3^5} \hat{z}$$

$$\langle 100 | \vec{r} | 21\pm 1 \rangle = \frac{2^7}{3^5} a (\mp \hat{x} - i \hat{y})$$

$$|200\rangle \rightarrow |100\rangle: \vec{p}^2 = 0$$

$$|210\rangle \rightarrow |100\rangle \quad |21\pm 1\rangle \rightarrow |100\rangle \quad \begin{aligned} |\vec{p}|^2 &= (p_x^2 + p_y^2 + p_z^2) \\ &= (ea)^2 \frac{2^{15}}{3^{10}} \end{aligned}$$

$$\omega_0 = \frac{E_2 - E_1}{\hbar} = \frac{1}{\hbar} \left(\frac{E_1 - E_1}{4} \right) = -\frac{3}{4\hbar} E_1$$

$$A = \frac{\omega_0^3 |p|^2}{3\pi\epsilon_0\hbar c^3} = -\frac{3^3 E_1^3}{2^6 \hbar^3} e^2 a^2 \frac{2^{15}}{3^{10}} \frac{1}{3\pi\epsilon_0\hbar c^3}$$

$$= \frac{2^9}{3^8} \left(-\frac{\hbar^2}{2ma^2} \right) \frac{E_1^2 e^2 a^2}{c^2} 4\alpha = \frac{2^{10} \alpha^4 m^2 c^4 \alpha}{\hbar m c^2}$$

$$\alpha = \frac{e^2}{4\pi\epsilon_0\hbar c}$$

$$A = \frac{Z^1}{3^3} \frac{\alpha^5 mc^2}{\hbar} = .15 \frac{\alpha^4 c}{\hbar}$$

$$= 6.27 \times 10^8 \text{ s}^{-1}$$

$$\Delta E = \hbar c A = 15 \alpha^3 Z^2 E_1 \\ = 10^7 E_1$$

$$\tau = \frac{1}{A} = 1.6 \times 10^{-9} \text{ s} \quad \text{for } n=2 \ell=1$$

$$\tau = \infty \quad \text{for } n=2 \ell=0$$

two photon emission $\tau = \frac{1}{10} \text{ s}$

Selection Rules

when is $\langle n'l'm' | \vec{r} | nlm \rangle = 0$?

$$[L_z, z] = 0$$

$$0 = \langle n'l'm' | L_z z - z L_z | nlm \rangle$$

$$= \hbar \langle n'l'm' | m' z - z m | nlm \rangle$$

$$= \hbar (m' - m) \langle n'l'm' | z | nlm \rangle$$

$$m' = m \quad \text{or} \quad \langle n'l'm' | z | nlm \rangle = 0$$

$$[L_z, x] = i\hbar y$$

$$\langle n'l'm' | [L_z, x] | nlm \rangle = \langle n'l'm' | L_z x - x L_z | nlm \rangle$$

$$i\hbar \langle n'l'm' | y | nlm \rangle = \hbar (m' - m) \langle n'l'm' | x | nlm \rangle$$

$$[L_z, y] = -i\hbar x$$

$$-i\hbar \langle n'l'm' | x | nlm \rangle = \hbar (m' - m) \langle n'l'm' | y | nlm \rangle$$

$$(m' - m)^2 = 1 \quad \text{or} \quad \langle n'l'm' | x | nlm \rangle = 0$$

$$\langle n'l'm' | y | nlm \rangle = 0$$

(1-1)

c' pole

so transitions allowed
on ℓ ; i.e. $\Delta m = \pm 1, 0$

$$[L^2, [L^2, \vec{r}]] = 2\hbar^2 (\vec{n} L^2 + L^2 \vec{n})$$

need $\Delta l = \pm 1$

photon carries spin 1

Prob. 9.20 Magnetic Resonance
spin $\frac{1}{2}$ particle gyromagnetic ratio γ

$$\vec{B} = B_{rf} \cos(\omega t) \hat{x} - B_{rf} \sin(\omega t) \hat{y} + B_0 \hat{z}$$

radio static

Larmor precess. $\omega_0 = \gamma B_0$

$$H = -\gamma \vec{B} \cdot \vec{S} = -\frac{\gamma \hbar}{2} (B_x \sigma_x + B_y \sigma_y + B_z \sigma_z)$$

$$= -\frac{\gamma \hbar}{2} \left(B_x \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix} + B_y \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix} + B_z \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix} \right)$$

$$= -\frac{\gamma \hbar}{2} \begin{pmatrix} B_0 & B_{rf} (\cos \omega t + i \sin \omega t) \\ B_{rf} (\cos \omega t - i \sin \omega t) & -B_0 \end{pmatrix}$$

$$= -\frac{\gamma \hbar}{2} \begin{pmatrix} B_0 & B_{rf} e^{i \omega t} \\ B_{rf} e^{-i \omega t} & -B_0 \end{pmatrix}$$

(To 2)

$$\text{at time } t: \quad X(t) = \begin{pmatrix} a(t) \\ b(t) \end{pmatrix}$$

$$it \dot{X} = HX$$

$$it \begin{pmatrix} \dot{a} \\ \dot{b} \end{pmatrix} = -\frac{\gamma t}{2} \begin{pmatrix} B_0 + B_{rf} e^{i\omega t} \\ B_{rf} e^{-i\omega t} - B_0 \end{pmatrix} \begin{pmatrix} a \\ b \end{pmatrix}$$

$$= -\frac{\gamma t}{2} \begin{pmatrix} B_0 a + B_{rf} e^{i\omega t} b \\ B_{rf} e^{-i\omega t} a - B_0 b \end{pmatrix}$$

$$\dot{a} = \frac{i\gamma t}{2} (B_0 a + B_{rf} e^{i\omega t} b) = \frac{i}{2} (\sqrt{2} e^{i\omega t} b + \omega_0 a)$$

$$\dot{b} = \cancel{\frac{i\gamma t}{2}} \left(\sqrt{2} e^{-i\omega t} a - \omega_0 b \right)$$

$$a(t) = \left(a_0 \cos(\omega' t/2) + \frac{i}{\omega'} (a_0(\omega_0 - \omega) + b_0 \sqrt{2}) \sin \frac{\omega' t}{2} e^{i\omega t/2} \right)$$

$$b(t) = \left(b_0 \cos(\omega' t/2) + \frac{i}{\omega'} (-b_0(\omega_0 - \omega) + a_0 \sqrt{2}) \sin \frac{\omega' t}{2} e^{-i\omega t/2} \right)$$

$$\omega' = \sqrt{(\omega - \omega_0)^2 + \sqrt{2}^2} \quad a(t=0) = a_0$$

check: $\ddot{a}(t) = \left(-\frac{\omega'}{2} a_0 \sin \frac{\omega' t}{2} + \frac{i}{2} (a_0(\omega_0 - \omega) + b_0 \sqrt{2}) \cos \frac{\omega' t}{2} \right) e^{i\omega t/2}$

$$+ i \frac{\omega}{2} a(t)$$

$$= \frac{i}{2} e^{i\omega t/2} \left(\omega a_0 \cos \frac{\omega' t}{2} + \frac{i\omega}{\omega'} (a_0(\omega_0 - \omega) + b_0 \sqrt{2}) \sin \frac{\omega' t}{2} \right.$$

$$\left. + i \omega' a_0 \sin \frac{\omega' t}{2} + (a_0(\omega_0 - \omega) + b_0 \sqrt{2}) \cos \frac{\omega' t}{2} \right)$$

(*)

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Compare with Schr. eq.

$$\ddot{a} = \frac{i}{2} (\sqrt{2} e^{i\omega t} b + w_0 a)$$

$$= \frac{i}{2} e^{i\omega t} \left(\sqrt{2} b_0 \cos \frac{\omega' t}{2} + i \frac{\sqrt{2}}{\omega'} (-b(w_0 - \omega) + a_0 \sqrt{2}) \sin \frac{\omega' t}{2} \right)$$

$$+ w_0 a_0 \cos \frac{\omega' t}{2} + i \frac{w_0}{\omega'} (a_0(w_0 - \omega) + b_0 \sqrt{2}) \sin \frac{\omega' t}{2}$$

 $\cos \frac{\omega' t}{2}$ terms match

from (*)

$$\sin \frac{\omega' t}{2} \text{ terms: } \left(i \frac{\omega}{\omega'} (a_0(w_0 - \omega) + b_0 \sqrt{2}) + i \omega' a_0 \right) \omega'$$

$$\stackrel{?}{=} i \sqrt{2} (-b_0(w_0 - \omega) + a_0 \sqrt{2}) + i w_0 (a_0(w_0 - \omega) + b_0 \sqrt{2})$$

$$a_0(\omega(w_0 - \omega) + \omega'^2 - \sqrt{2}^2 - w_0^2 + w_0 \omega) \stackrel{?}{=} b_0(-\sqrt{2}(w_0 - \omega) + w_0 \sqrt{2} - w_0 \omega)$$

$$a_0(\omega(w_0 - \omega) + (w_0 - \omega)^2 - w_0^2 + w_0 \omega) \stackrel{?}{=} 0$$

$$a_0(2\omega w_0 - \omega^2 + w_0^2 - 2w_0 \omega + \omega^2 - w_0^2) = 0$$

solution works